

Lecture#2: Energy loss of electrons in matter and Photon interactions

2 main differences:

- Projectile identical to target electron
- Bremsstrahlung process important since m_e small (on $+Ze$ & e^-)
- Moller scattering, Bhabha scattering, e^+e^- annihilation

$$S_{\text{coll}} = K \cdot \frac{1}{2} \left\{ \ln \frac{m\beta^2 c^2 E}{2I^2(1-\beta^2)} - \ln 2 [2(1-\beta^2)^{1/2} - 1 + \beta^2] \right. \\ \left. + (1-\beta^2) + \frac{1}{8} [1 - (1-\beta^2)^{1/2}]^2 \right\}$$

$$S_{\text{rad}} = NEZ(Z+1)e^4 \{ 4 \ln (2E/mc^2) - 4/3 \} / [137m^2c^4]$$

$$S_{\text{tot}} = S_{\text{coll}} + S_{\text{rad}}$$

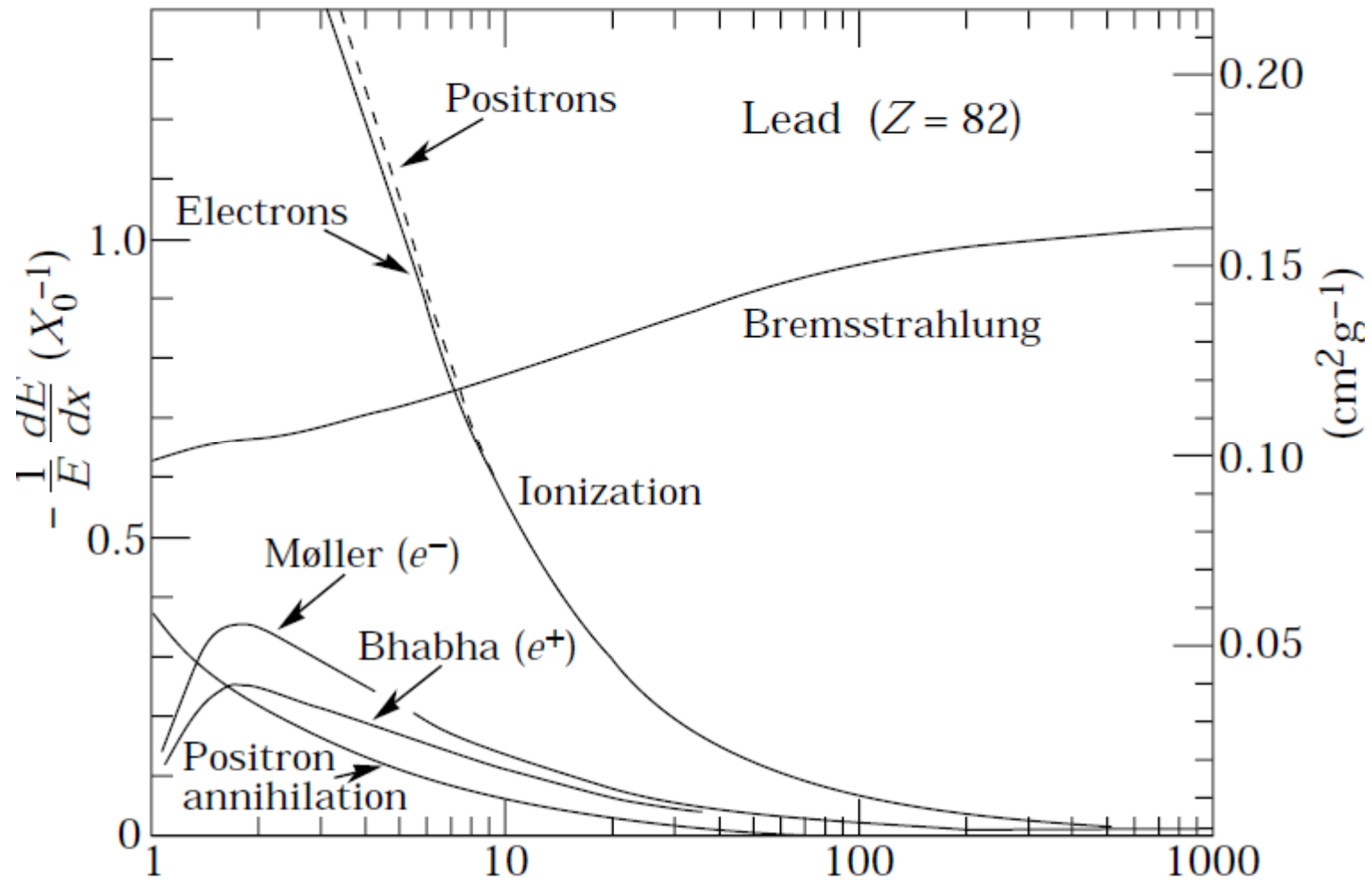
$$E_c = 800 \text{ MeV}/(Z+1.2) \text{ where } S_{\text{coll}} \approx S_{\text{rad}}$$

Radiation length X_0 defined as

- electron loses energy to $1/e$ of its initial value by bremsstrahlung
- $7/9$ of γ -attenuation length (or m.f.p.)

$$X_0 = 716.4 \text{ A} / [Z(Z+1) \ln (287/\sqrt{Z})] \quad \text{Dahl}$$

Energy loss of e^- and e^+ in Pb



$X_0 (\text{Pb}) = 6.37 \text{ gm/cm}^2$

$E (\text{MeV})$

Backscattering of low energy electrons

$$f_{\text{backsc}} = N_{\text{e}}^{\text{backsc}} / N_{\text{inc}}$$

About 0.14, 0.3, 0.5 for Si, Ge, Au respectively at low energy
and at angle of incidence $\theta_{\text{inc}} = 0^\circ$ (normal to surface)

Drops off at higher energies (about half its value at ~ 6 MeV)

Increases with θ_{inc}

Interaction of γ -rays with matter

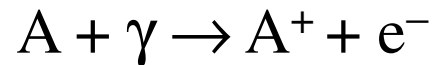
3 main processes:

1. Photoelectric effect ($\sigma_{\text{pe}} \propto Z^5 / E^{3.5}$)

2. Compton scattering ($\sigma_{\text{C}} \propto Z / E$)

3. Pair production $\sigma_{\text{pair}} \propto Z^2 \ln(E_{\gamma} / m_e c^2)$

Photoelectric effect



Threshold energy : Binding energy of electron in
atom/molecule/solid...

$E_{pe} = E_\gamma - BE \Rightarrow$ Full energy deposited in detector (K- or L-Xray
escape?)

$$\sigma_{pe} = 4\alpha^4 \sqrt{2} Z^5 \sigma_0 (mc^2/E_\gamma)^{7/2} \text{ where } \sigma_0 = 8\pi r_e^2 / 3 \text{ for } E_\gamma > BE_K$$

For E_γ higher than, but close to, K-edge

$$\sigma_{pe} = (6.3 \times 10^{-18} / Z^2) (BE_K/E_\gamma)^{8/3}$$

Compton scattering

For a photon scattering off a free electron at rest

$$\gamma + e^- \rightarrow \gamma + e^-$$

$$E_\gamma = E_{\gamma 0} / [1 + E_\gamma / m_e c^2 (1 - \cos \theta)]$$

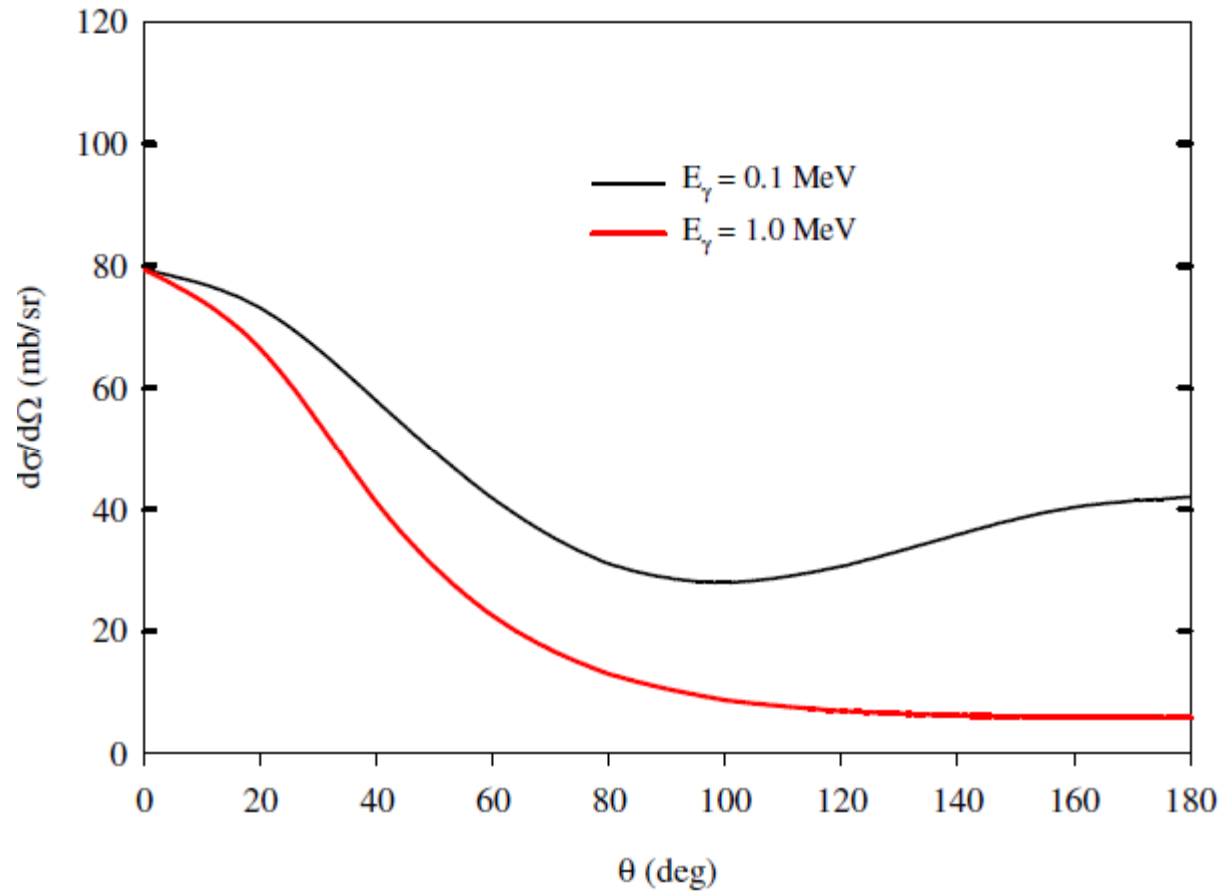
For a bound electron there will be smearing of this energy and will increase the more bound the electron (*used to measure momentum distribution in solids!*)

Limiting cases: 1. If $E_{\gamma 0} \ll m_e c^2$, $E_\gamma \approx E_{\gamma 0}$ at all angles

2. If $E_{\gamma 0} \gg m_e c^2$, $E_\gamma \approx E_{\gamma 0}$ at near 0° , and $\frac{1}{2} m_e c^2$ for $\theta \approx 180^\circ$

Here $E_e \approx E_{\gamma 0} - \frac{1}{2} m_e c^2$, $dE_e/d\theta \approx 0 \Rightarrow$ *quasi-monoenergetic* γ !

Angular distribution of Compton scattered gamma-rays



$$d\sigma/d\Omega = \frac{1}{2} r_e^2 \left\{ \frac{1}{[1 + \epsilon(1 - \cos \theta)]^2} \right\} \left\{ 1 + \cos^2 \theta + \frac{\epsilon^2 (1 - \cos \theta)^2}{1 + \epsilon(1 - \cos \theta)} \right\}$$

where $\epsilon = E_\gamma/(m_e c^2)$, $r_e = e^2/m_e c^2 = 2.818$ fm is classical electron radius

➤ Compton scattering leads to continuous energy

deposit upto Compton edge $E_{\gamma}^{\text{CE}} = E_{\gamma 0} - m_e c^2/2$

➤ In large HPGe detectors Compton followed by photoabsorption leads to full energy peak

➤ Compton scattering cross section also depends on photon and electron spin polarisations (formula above has averaged over both spins and summed over final ones)

An application of Compton effect: Laser backscattered photons

Suppose you shoot a laser beam (in visible spectrum, for example) at an incoming relativistic electron beam

In rest frame of electron, photons are Doppler shifted by factor

$$\gamma(1 - \beta \cos \theta_{e\gamma}) \approx 2\gamma \text{ for } \theta_{e\gamma} = 180^\circ$$

When it backscatters photon, its Doppler shifted to still higher energies in lab, by another factor of $\sim 2\gamma$

\Rightarrow Overall increase in energy by $\sim 4\gamma^2$

For $\gamma \sim 10^3$ ($E_e = 500 \text{ MeV}$) $E_\gamma \sim 10 \text{ MeV}$

and $\gamma \sim 10^4$ ($E_e = 5 \text{ GeV}$) $E_\gamma \sim 1 \text{ GeV!}$

Pair production

Photons with energy $> 2m_e c^2$ can produce $e^+ e^-$ pairs

However **energy-momentum** conservation requires this to occur near an electric charge such as a *nucleus* or *electron*

$$\gamma + Ze \rightarrow e^+ + e^- + Ze$$

$$\gamma + e^- \rightarrow e^+ + e^- + e^-$$

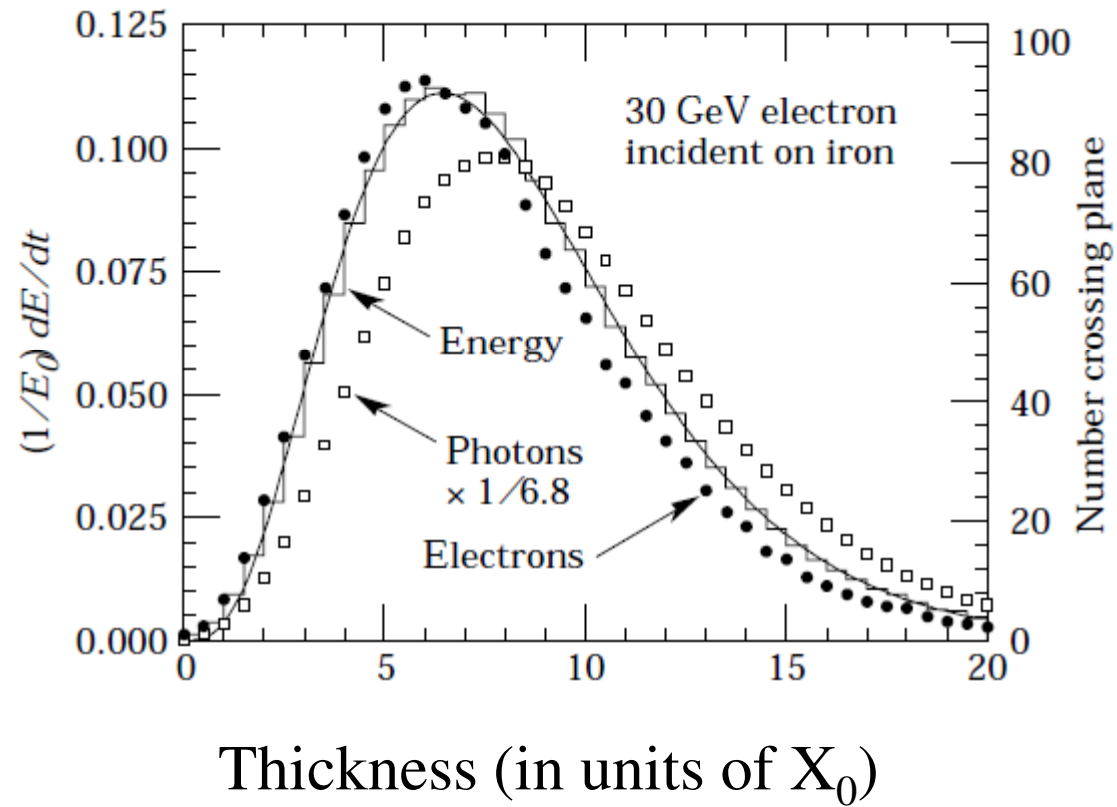
$$\sigma_{\text{pair}} = 4Z^2 \alpha r_e^2 \left[\frac{7}{9} \{ \ln (2E_\gamma / m_e c^2) - f(Z) \} - 109/54 \right]$$

Electromagnetic shower for $E_\gamma \gg 1.02 \text{ MeV}$

Electron-positron pair produced \rightarrow bremsstrahlung $\rightarrow e^+ + e^-$ pair and process repeats (Bhabha – Heitler). First discussed in the context of cosmic ray showers.

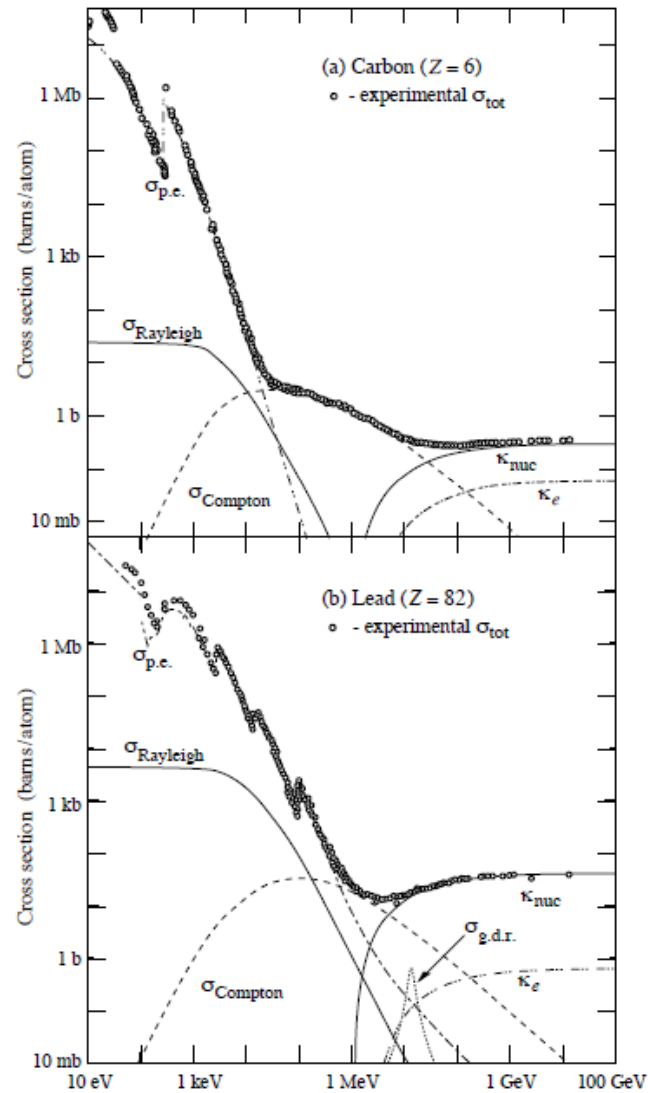
In detector max. energy deposition via photoelectric and pair production processes. Compton process leads to only partial energy deposit. In EM shower 511 keV annihilation photons and bremsstrahlung can escape leading to tailing on the lower energy side of full energy peak

EGS4 simulation of 30 GeV electron in Iron

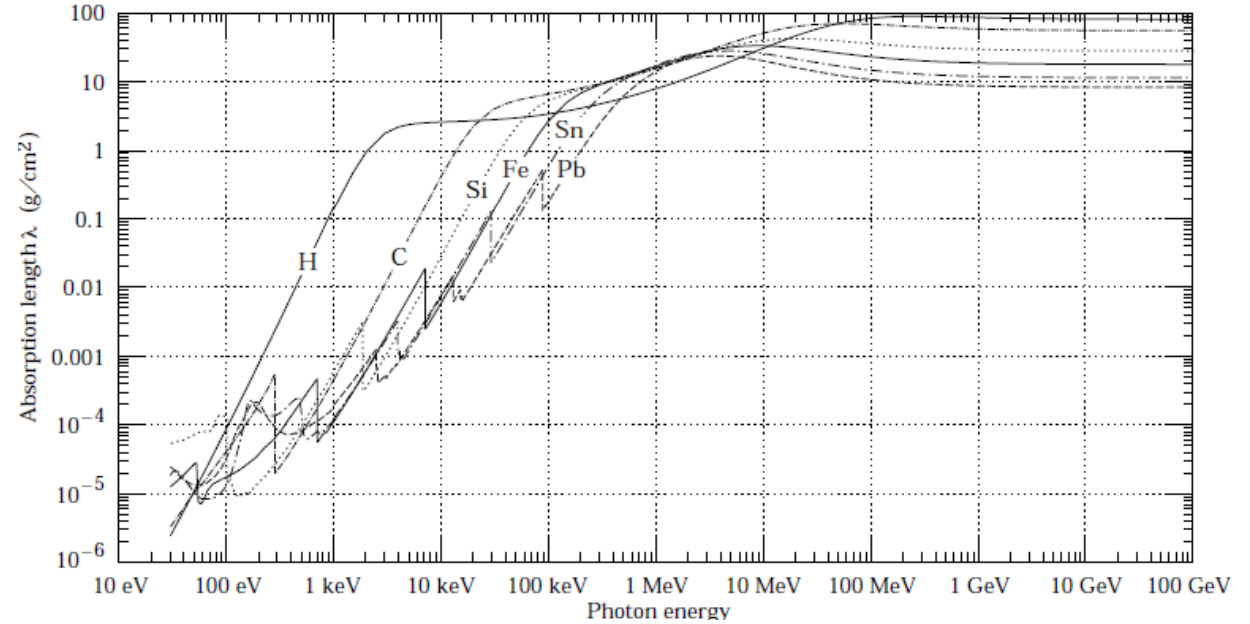


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Photon interaction with lead



Photon energy



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