# Flavour physics – Lecture 2

Jim Libby (IITM) VIII<sup>th</sup> SERC school on EHEP VECC Kolkata – 30<sup>th</sup> June to 2<sup>nd</sup> July

# Outline

- CKM again and its predictions
  - Charm, bottom and top discovery
- Mixing
- Types of CP violation
- Kaon physics
- Charm physics

# CKM matrix

- Extension of Cabibbo's two by two mixing matrix
- Unitary complex matrix
   18 parameters
- Unitary constraints (VV+=I)
  - 9 parameters
- Quark fields can absorb five unobservable phases
  - 4 parameters
  - a 3 mixing angle and 1 phase

$$\begin{pmatrix} u & c & t \end{pmatrix} \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

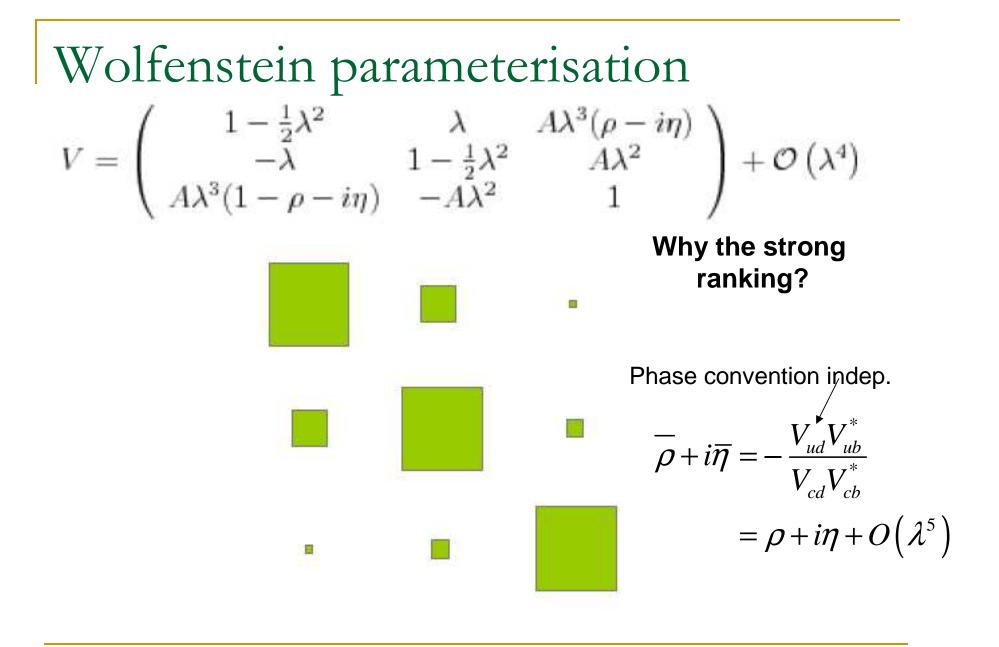
#### Responsible for CP violation

# CKM matrix parameterisations

3 mixing angles (θ<sub>12</sub>, θ<sub>13</sub>, θ<sub>23</sub>)+ one phase (δ)
 PDG (Chau and Keung) - s<sub>ij</sub>=sin θ<sub>ij</sub> and c<sub>ij</sub>=cos θ<sub>ij</sub>

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

- $\Box S_{12} \sim 0.2$   $S_{23} \sim 0.04$   $S_{13} \sim 0.004$
- However, Wolfenstein parameterisation that exploits this hierarchy and expands matrix in terms of λ=s<sub>12</sub>
  - then  $s_{23} = A\lambda^2$  and  $s_{13}e^{i\delta} = A\lambda^3(\rho + i\eta)$



Unitarity  
$$V^{\dagger}V = VV^{\dagger} = I$$

$$\begin{aligned} |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 &= 1\\ |V_{cd}|^2 + |V_{cs}|^2 + |V_{cb}|^2 &= 1\\ |V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2 &= 1 \end{aligned}$$

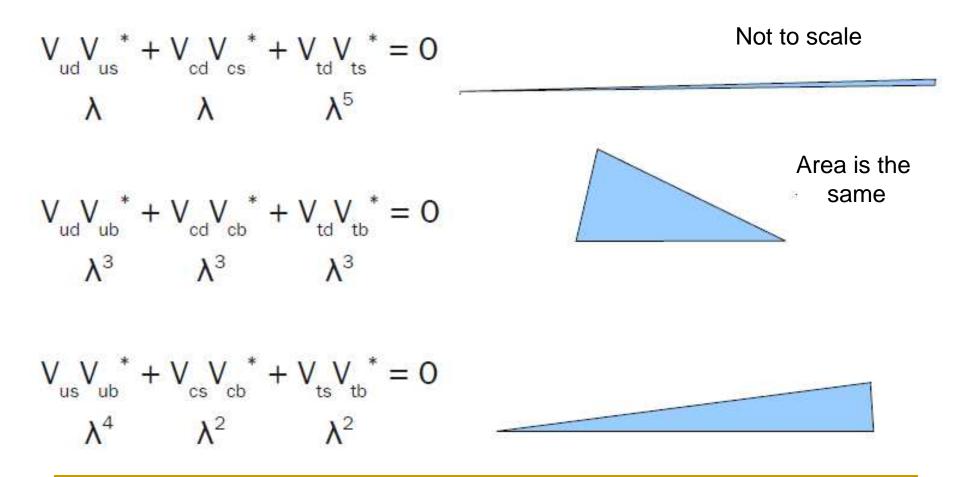
$$\begin{split} |V_{ud}|^2 + |V_{cd}|^2 + |V_{td}|^2 &= 1 \\ |V_{us}|^2 + |V_{cs}|^2 + |V_{ts}|^2 &= 1 \\ |V_{ub}|^2 + |V_{cb}|^2 + |V_{tb}|^2 &= 1 \end{split}$$

$$V_{ud}V_{us}^{*} + V_{cd}V_{cs}^{*} + V_{td}V_{ts}^{*} = 0$$
  
$$V_{ud}V_{ub}^{*} + V_{cd}V_{cb}^{*} + V_{td}V_{tb}^{*} = 0$$
  
$$V_{us}V_{ub}^{*} + V_{cs}V_{cb}^{*} + V_{ts}V_{tb}^{*} = 0$$

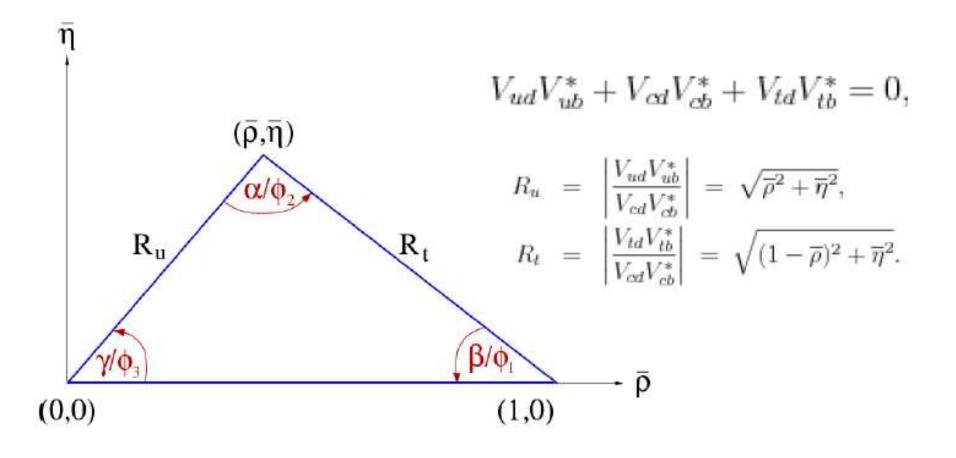
$$V_{ud}V_{cd}^{*} + V_{us}V_{cs}^{*} + V_{ub}V_{cb}^{*} = 0$$
  
$$V_{ud}V_{td}^{*} + V_{us}V_{ts}^{*} + V_{ub}V_{tb}^{*} = 0$$
  
$$V_{cd}V_{td}^{*} + V_{cs}V_{ts}^{*} + V_{cb}V_{tb}^{*} = 0$$

# Unitarity triangles

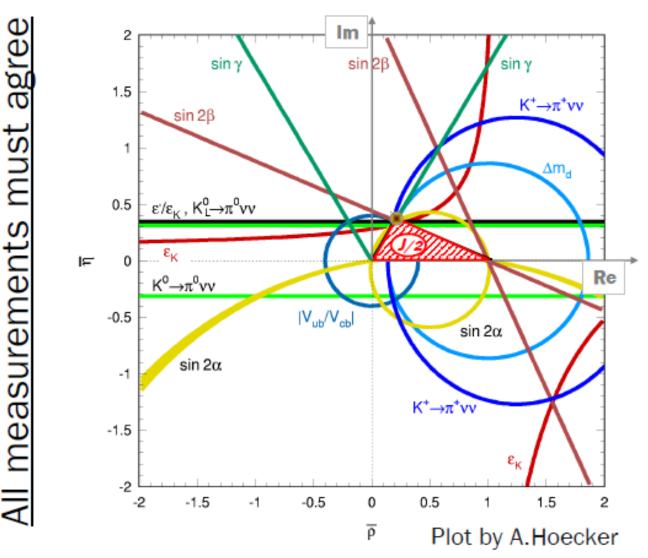
Null relations can be expressed as triangles in complex plane



# The unitarity triangle



#### Predictive nature of KM mechanism



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#### But the first KM (and GIM) predictions

#### More quarks

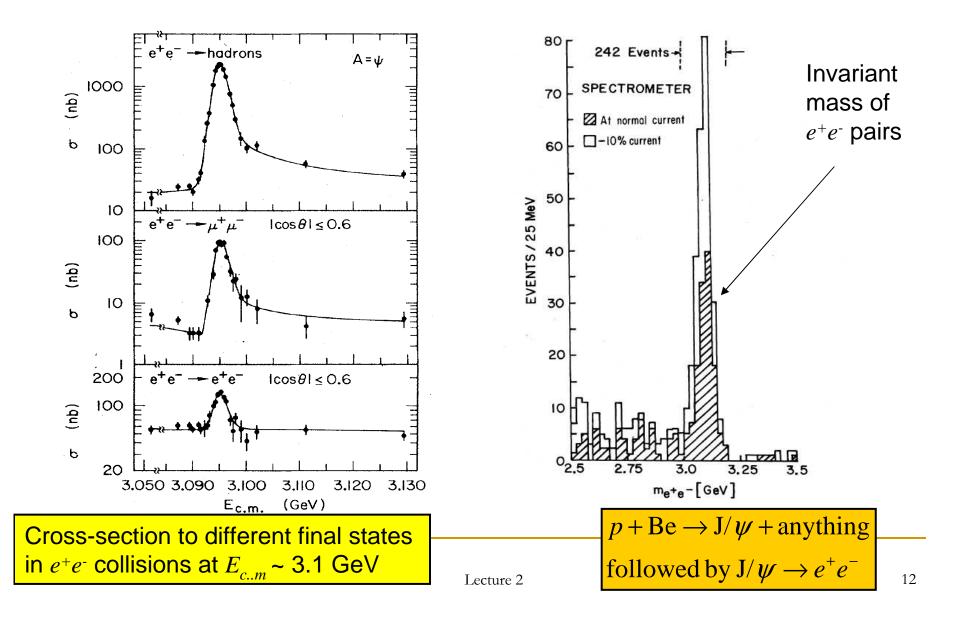
- Discovery of charm then beauty (and finally top)
- Worth spending a little time as it highlights a couple of features of experiments
- Today principle interest now in onia is QCD
  - Exotic X, Y, Z states found at Babar and Belle
    - New states D-molecules, tetraquarks etc
  - Some states only just observed

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# Discovery of Charmonium

- The charm quark had been proposed in 1970 by Glashow, Illiopoulos and Maiani to explain the suppression of transitions involving flavour changing neutral currents such as  $K^0 \rightarrow \mu\mu$  which  $\Delta S = 1$ 
  - □ Flavour changing charged currents had been observed since  $\beta$  decay which is  $u \rightarrow de^+ v_e$
- The first direct evidence for the charm quark came in 1974 from two different experiments:
  - $e^+e^-$  collisions at a centre of mass energy ~3.0 GeV/ $c^2$ 
    - SPEAR collider at Stanford Linear Accelerator Center
  - $\square$  p+Be fixed target experiment with 28 GeV protons
    - Brookhaven alternating-gradient synchrotron

# Discovery of the $J/\psi$



#### More charmonium

- An additional resonance was soon observed at SLAC with a mass of 3686 MeV and was called the ψ' or ψ(2S)
- Similar arguments to those used for the J/ψ led to the same assignment of quantum numbers: J<sup>PC</sup>=1<sup>--</sup>, I=0
- The ψ' is narrow Γ=0.28 MeV but there were new decay modes observed:
  - $\Box \quad \psi' \rightarrow J/\psi + 2\pi \ (BR = 50\%)$
  - $\Box \quad \psi' \rightarrow \chi + \gamma \; (BR = 24\%)$

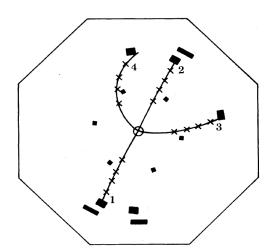
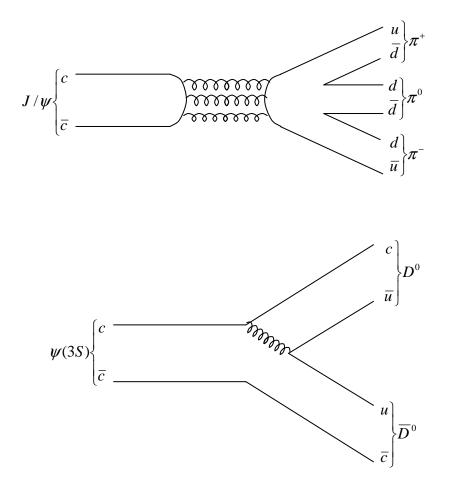


FIG. 3. An example of the decay  $\psi(3684) \rightarrow \pi^+ + \pi^-$ + $\psi(3095)$ , where  $\psi(3095) \rightarrow e^+ + e^-$ , from an off-line reconstruction of the data. The event is seen in the *x*-*y* projection where *z* is the beam (and magnetic field) direction. Also shown are the trigger and shower counters which detected the tracks. Tracks 3 and 4 are the slow pions and tracks 1 and 2 are the two leptons from  $\psi(3095)$  decay.

• There were a further 4 states observed labelled  $\psi(3S)$ ,  $\psi(4S)$  and  $\psi(5S)$  – but their widths were much greater between 24 and 78 MeV

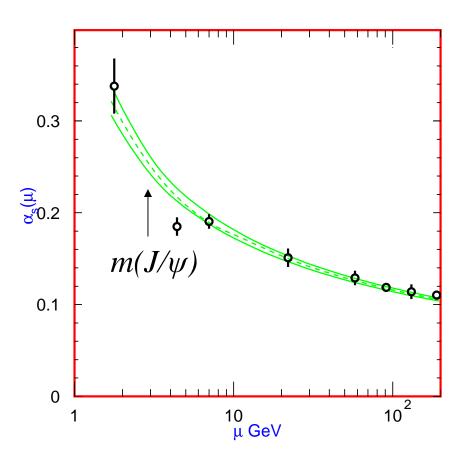
# Decays of the charmonium

- Allowed hadronic final states governed by kinematics
   m(J/ψ) < m(ψ') < 2m(D) < m(ψ(3S))</li>
- Strong decay of the  $J/\psi$  is via the annihilation of the charm and anticharm quarks to 3 gluons
  - An odd number of gluons is required to conserve C
    - Gluon and  $J/\psi$  both have C=-1
  - A single gluon is not colourless therefore is not colour charge conserving
  - $\square \quad \alpha_s^{\ 6}$  dependence
- Strong decay of  $\psi(3S)$  is a single gluon radiated to form *D* meson pair
  - $\square \quad \alpha_s^2$  dependence



# Running of the strong coupling

- Additional factor comes the strength of the strong force varying rapidly as a function of the momentum transfer of the interaction µ
- The coupling is much stronger for a  $\mu \sim m(u/d)$  than for  $\mu \sim m(J/\psi)$

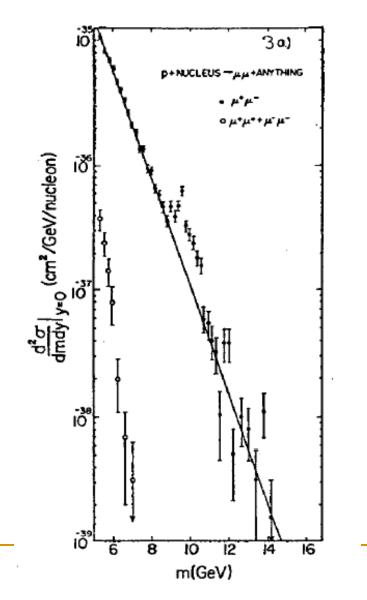


#### Bottomonium

- In 1977 there were a similar set of resonances observed in the centre of mass region 9.5-10.5 GeV
- Observed in interactions of proton beam on fixed targets:

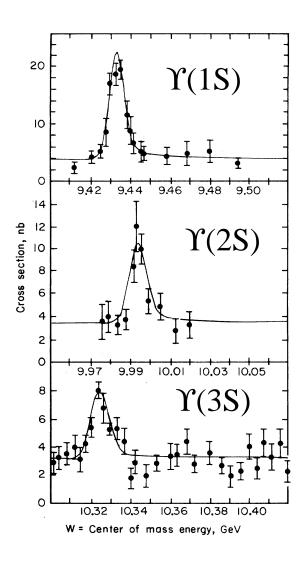
 $p + \text{Be}, \text{Cu}, \text{Pt} \rightarrow \mu^+ + \mu^- + \text{anything}$ 

 Resolution of 500 MeV but width of resonance 1.2 GeV suggested it might be several resonances



#### Bottomonium

- Resonances were then observed in e<sup>+</sup>e<sup>-</sup> collisions and they were resolved into three narrow states Υ(1S), Υ(2S) and Υ(3S)
- A broader resonance Y(4S) that decays into B meson pairs was also found
- Associated with bound states of b quarks and antiquarks with Y(4S) the first resonance to be heavier than 2 B mesons



# Mixing

- For neutral particles with opposite flavour quantum numbers (i.e. K0 and K0bar) CPT theorem requires:
  - Equal masses
  - Equal lifetimes
- However, mixed states K<sub>S</sub> and K<sub>L</sub> can have different masses and widths

# Mixing formalism

# Arbitrary mixture of $B^0$ and $\overline{B}^0$

$$p\left|B^{0}\right\rangle + q\left|\overline{B}^{0}\right\rangle$$

Evolution governed by time-dependent Schrodinger equation

$$i\frac{\partial}{\partial t}\binom{p}{q} = H\binom{p}{q} = \left(M - \frac{i}{2}\Gamma\right)\binom{p}{q}$$

H is Hamiltonian

*M* and  $\Gamma$  are 2×2 Hermitian matrices

Encapsulate the perturbative weak effects that lead to mixing CPT requires the diagonal elements of M and  $\Gamma$  are equal

# Mixing formalism

Heavy and light mass eigenstates

Eigenvalue equation gives

$$\begin{aligned} |B_L\rangle &= p|B\rangle + q|\overline{B}\rangle \\ |B_H\rangle &= p|B\rangle - q|\overline{B}\rangle, \\ \frac{q}{p} &= \sqrt{\frac{M_{12}^* - \frac{i}{2}\Gamma_{12}^*}{M_{12} - \frac{i}{2}\Gamma_{12}}}, \end{aligned}$$

$$\lambda_{H,L} = m_{H,L} - \frac{i}{2} \Gamma_{H,L},$$

Where  $m_H, m_L = M \pm Re \sqrt{|M_{12}|^2 - \frac{|\Gamma_{12}|^2}{4} - iRe(M_{12}\Gamma_{12}^*)} \equiv M \pm \Delta m/2,$  $\Gamma_H, \Gamma_L = \Gamma \pm 2Im \sqrt{|M_{12}|^2 - \frac{|\Gamma_{12}|^2}{4} - iRe(M_{12}\Gamma_{12}^*)} \equiv \Gamma \pm \Delta \Gamma/2,$ 

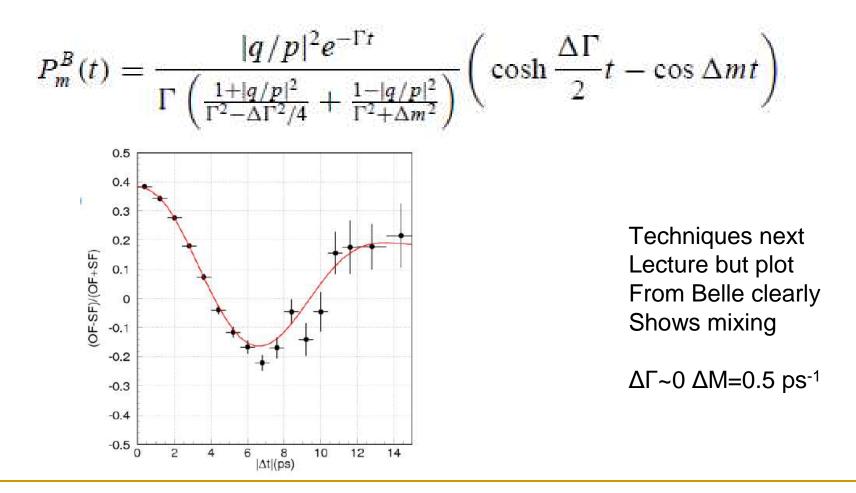
# Mixing formalism

Mass and width difference related To off diagonal elements

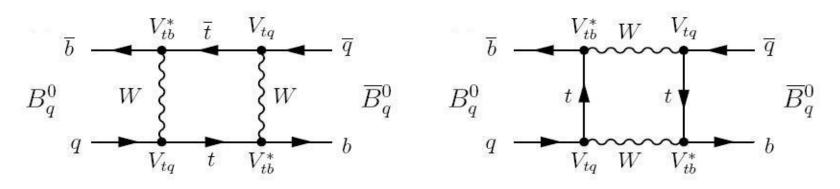
$$\Delta m^2 - \frac{\Delta \Gamma^2}{4} = 4 \left( |M_{12}|^2 - \frac{|\Gamma_{12}|^2}{4} \right)$$
$$\Delta m \Delta \Gamma = 4 \operatorname{Re}(M_{12}\Gamma_{12}^*).$$

$$\begin{split} |B(t)\rangle &= \frac{1}{2p}(|B_L(t)\rangle + |B_H(t)\rangle) & \text{Can express time-evolution} \\ &= \frac{1}{2p}(|B_L(t)\rangle + |B_H(t)\rangle) & \text{can express time-evolution} \\ &= \frac{1}{2}e^{-iMt}e^{-\frac{\Gamma}{2}t} \left( \left( e^{\frac{\Delta\Gamma}{4}t}e^{i\frac{\Delta m}{2}t} + e^{-\frac{\Delta\Gamma}{4}t}e^{-i\frac{\Delta m}{2}t} \right) |B\rangle \right. \\ &+ \frac{q}{p} \left( e^{\frac{\Delta\Gamma}{4}t}e^{i\frac{\Delta m}{2}t} - e^{-\frac{\Delta\Gamma}{4}t}e^{-i\frac{\Delta m}{2}t} \right) |\overline{B}\rangle \right). \\ &\text{Similar for Bbar} \end{split}$$

Mixing formalismProbability to mix



# Mixing



Feynman box diagrams for B mixing.

Weak box diagrams drive mixing

Sensitive to CKM elements Vtd (Vts) for B (Bs) mixing

More later on relating to unitarity triangle

# Types of CP violation

• Consider decay of neutral particle to a CP eigenstate  $\Delta m$ 

$$\frac{|\frac{q}{p}| \neq 1}{\frac{1}{4}}$$

 $\left|\frac{A}{A}\right| \neq 1$ 

 $\Im\left(\frac{q}{n}\frac{\overline{A}}{A}\right) \neq 0$ 

**CP violation** in mixing

**CP violation in decay (direct CPV)** 

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CP violation in interference between mixing and decay

Lecture 2

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# CP violation kaons

- CP violation in neutral kaon system due to |p/q|=(1+ε)/(1-ε)≠1 and possibility of in decay (ε')
- Observables studied

$$\begin{split} A_L &= \frac{\Gamma(K_L^0 \to \pi^- \ell^+ \nu) - \Gamma(K_L^0 \to \pi^+ \ell^- \nu)}{\Gamma(K_L^0 \to \pi^- \ell^+ \nu) + \Gamma(K_L^0 \to \pi^+ \ell^- \nu)} = 2 \operatorname{Re} \epsilon / (1 + |\epsilon|^2) \approx 2 \operatorname{Re} \epsilon \\ \eta_{+-} &= A(K_L^0 \to \pi^- \pi^-) / A(K_S^0 \to \pi^+ \pi^-) \qquad |\epsilon| = (2.228 \pm 0.011) \times 10^{-3} , \\ &= |\eta_{+-}| \ e^{i\phi_{+-}} = \epsilon + \epsilon' \qquad \text{PDG values} \\ \eta_{00} &= A(K_L^0 \to \pi^0 \pi^0) / A(K_S^0 \to \pi^0 \pi^0) \\ &= |\eta_{00}| \ e^{i\phi_{00}} = \epsilon - 2\epsilon' \qquad \epsilon' / \epsilon = (1.65 \pm 0.26) \times 10^{-3} \end{split}$$

# Direct CP violation

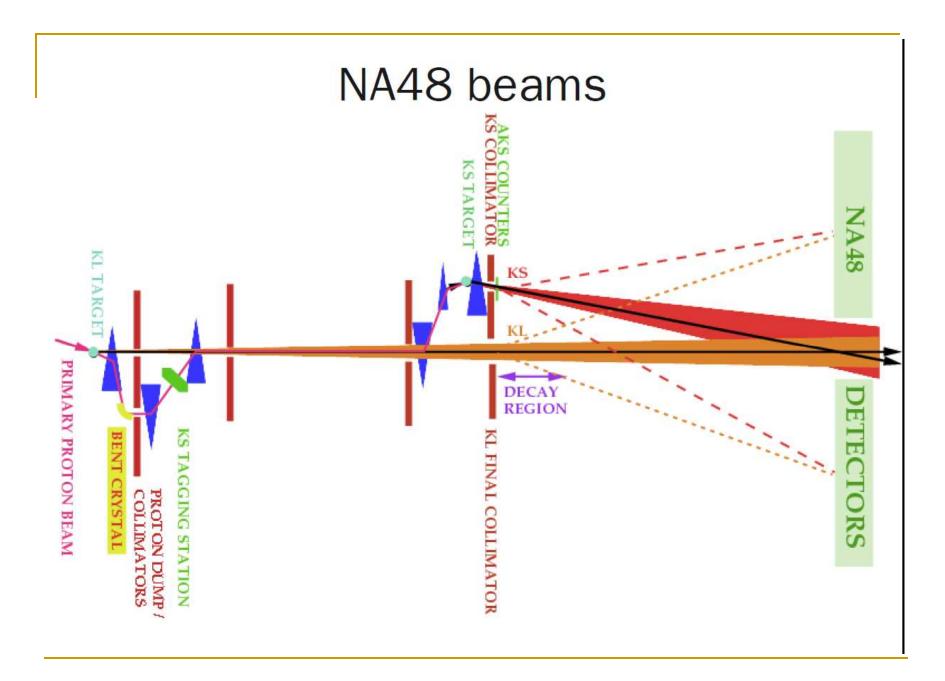
- New superweak force not KM gives CP (Wolfenstein)
- There would only be CPV in K mixing and not in other system
- Long search in kaon physics to sort this out
   NA48 and KTeV to measure
  - if CPV in  $K_{L} \rightarrow \pi^{+}\pi^{-} \neq$  CPV in  $K_{L} \rightarrow \pi^{0}\pi^{0} \leftrightarrow$  CPV cannot be in mixing only

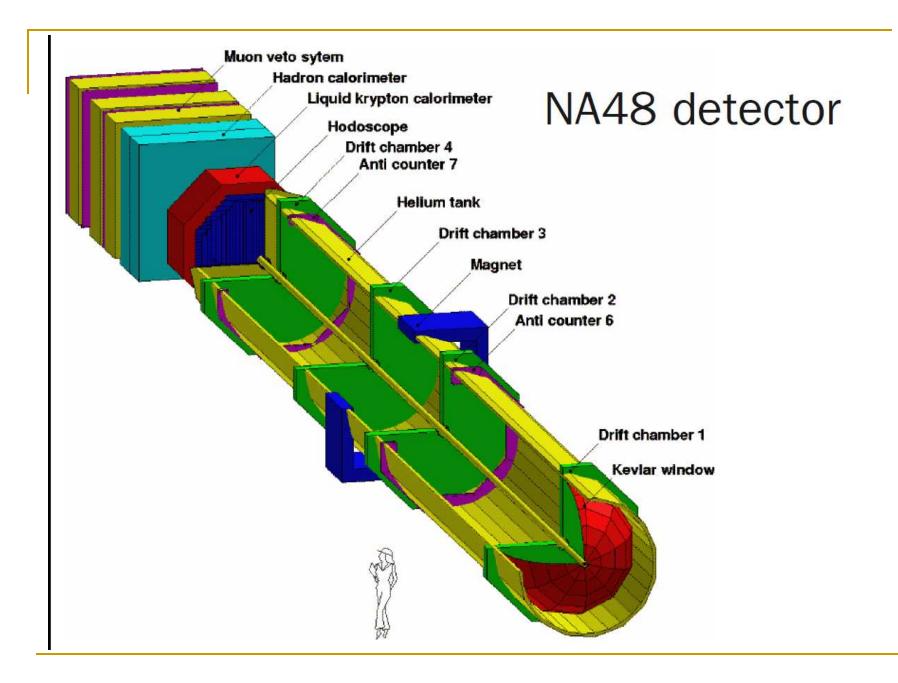
$$\mathsf{R} = (\mathsf{K}_{_{\mathsf{L}}} \to \pi^{_{0}}\pi^{_{0}})/(\mathsf{K}_{_{\mathsf{S}}} \to \pi^{_{0}}\pi^{_{0}})/(\mathsf{K}_{_{\mathsf{L}}} \to \pi^{^{+}}\pi^{\bar{}})/(\mathsf{K}_{_{\mathsf{S}}} \to \pi^{^{+}}\pi^{\bar{}})$$

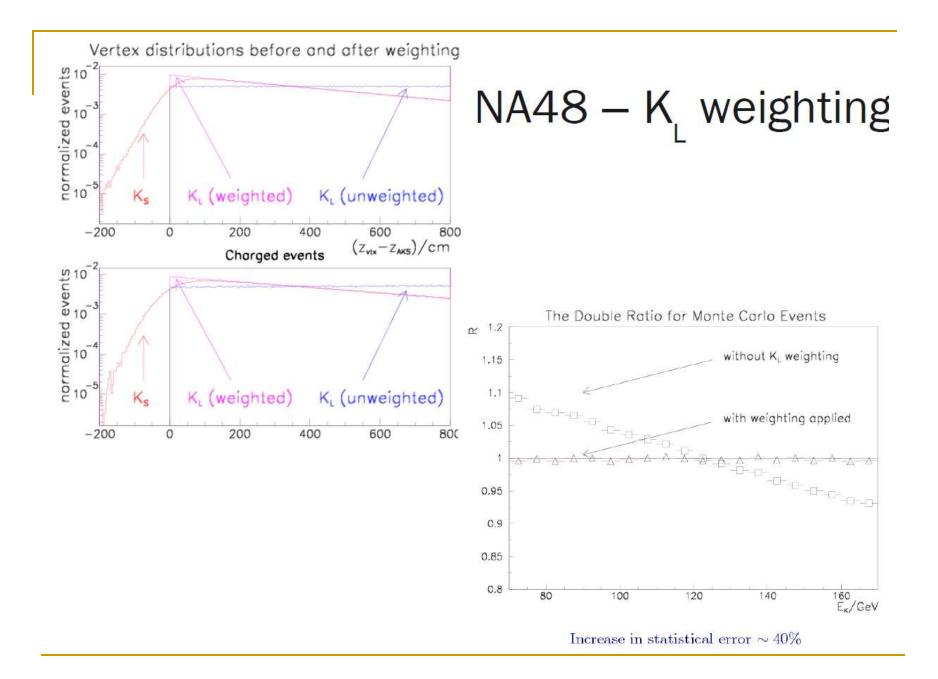
 $R = 1 - 6 \operatorname{Re}(\epsilon'/\epsilon)$ 

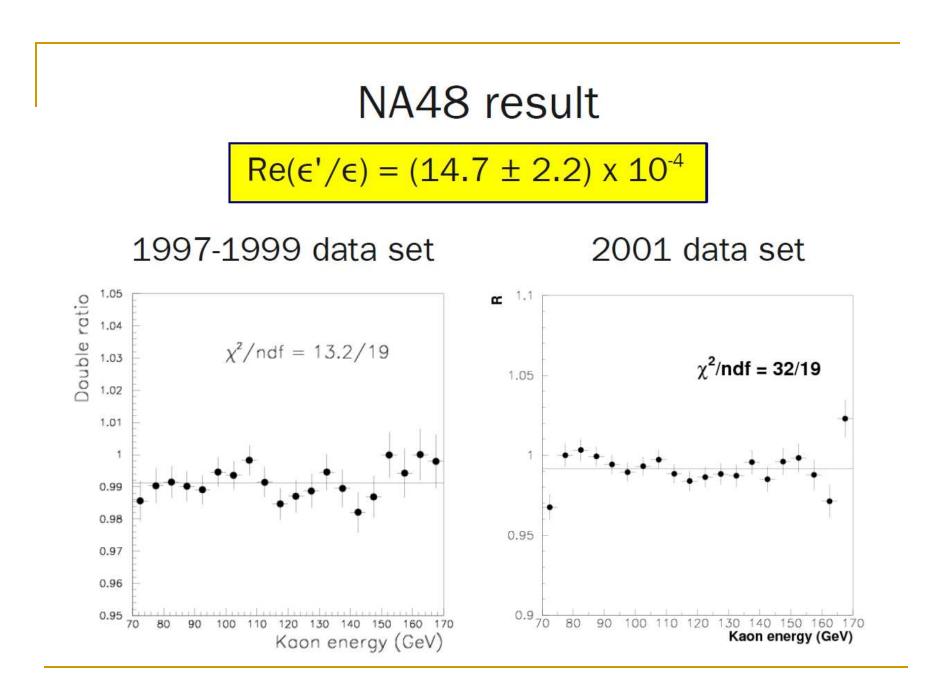
#### NA48 technique

- Precise measurement need control of systematic errors
  - use simultaneous  $K_{\!_{\rm I}}$  and  $K_{\!_{\rm S}}$  beams
  - take data in all four modes together
  - make acceptance as similar as possible
  - perform analysis in bins on kaon energy
    - correct for differences in K<sub>1</sub> and K<sub>s</sub> energy spectra
  - weight  $\rm K_{_{\rm I}}$  events according to  $\rm K_{_{\rm S}}$  decay distribution

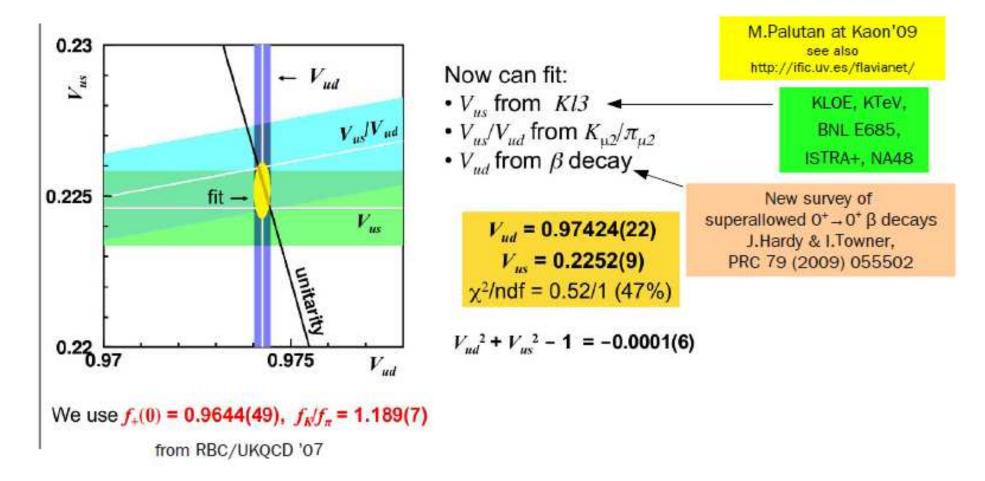








# Current and future kaon physics Precision measurements of |V<sub>us</sub>|



#### Future kaon physics

- "Holy grail" of kaon physics : rare decays  $K \rightarrow \pi v v$ 
  - Allow precise measurements of CKM matrix parameters
  - Sensitive to new physics effects
  - Extraordinarily challenging experimentally
  - $K^+ \rightarrow \pi^+ VV$ 
    - Current measurement (1.7 ± 1.1) 10<sup>-10</sup> BNL-E949 (expectation 0.9  $10^{-10}$ )
    - NA62 experiment (CERN) will observe O(100) SM events
  - $K^0 \rightarrow \pi^0 V \underline{V}$
- Current UL < 6.7  $10^{-8}$  KEK-E391 (expectation 2.5  $10^{-11}$ )
  - KOTO experiment (JPARC) will reach SM sensitivity

- Height of UT

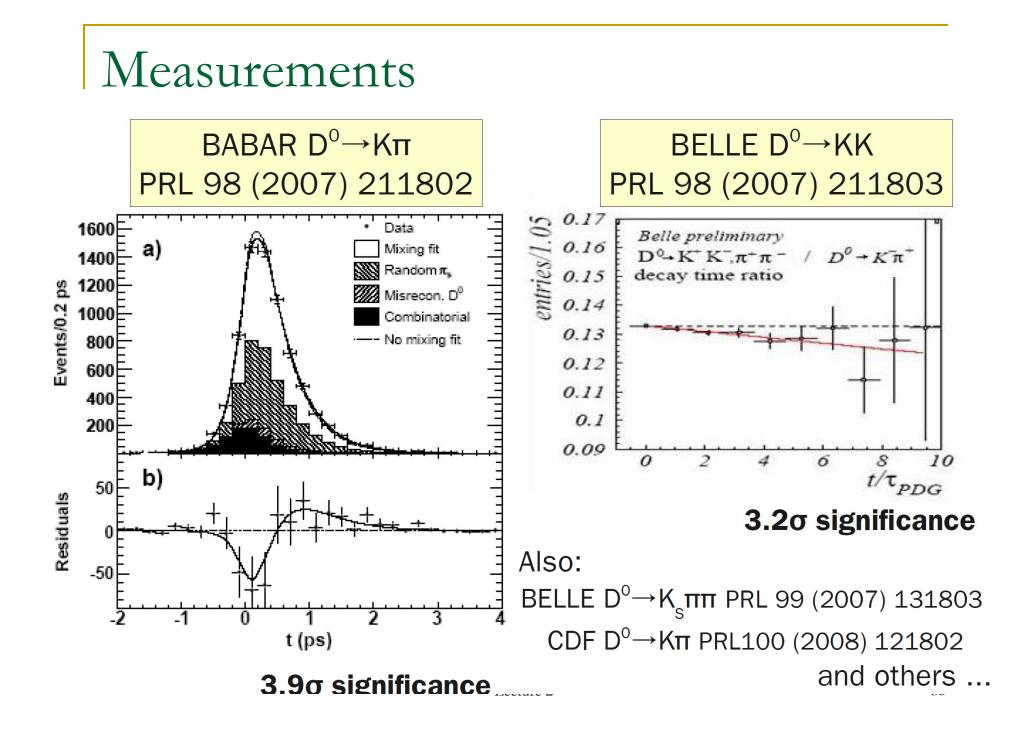
to. *n* 

# Charm physics

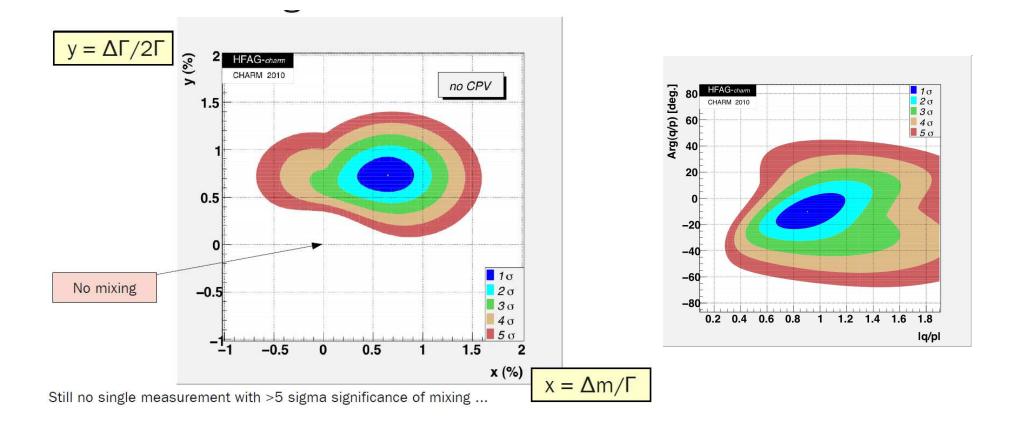
- Mixing has down type quarks in the box
- Only place this can be studied as top quark does not hadronise
- Also, SM CPV very small

Any CPV is new physics

- B-factories have produced lots of Ds and observed D-mixing
  - May yet see CPV
  - Big target of flavour physics at the LHC



#### Combination of measurements



#### Mixing confirmed but no CPV in mixing

# Lecture 2 - summary

CKM matrix and the KM mechanism

- Prediction of c and b their discovery
- Neutral meson mixing
- Different types of CPV
  - Kaon CPV yes
  - Charm CPV no