

# Lecture I

---

## Detector Instrumentation

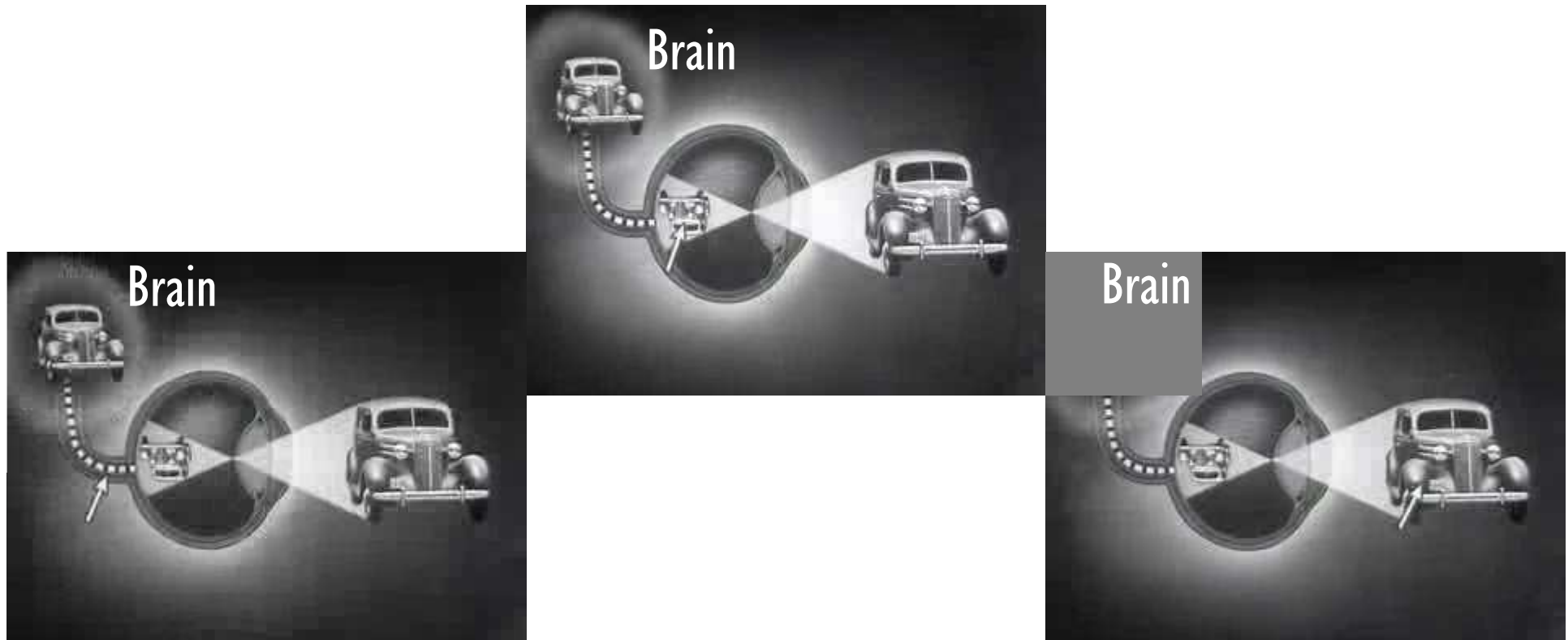
### Introduction

*Reference for this lecture series:*

*William Leo, Techniques of Nuclear and Particle Physics*

# How do you see ?

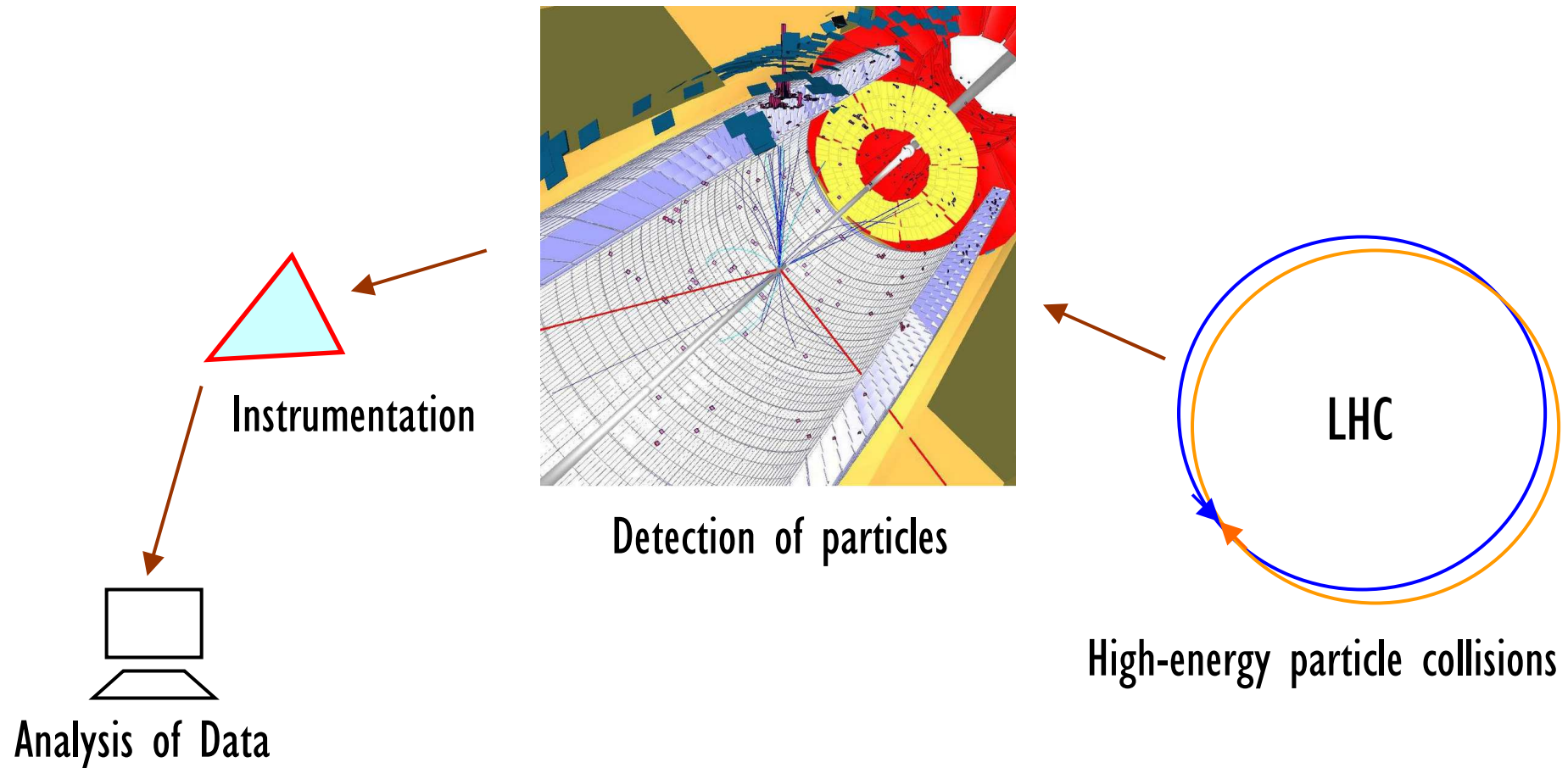
---



Instrumentation ← Detection ← Phenomenon

# Detector instrumentation

---



# Why do you want to study fundamental particles ?

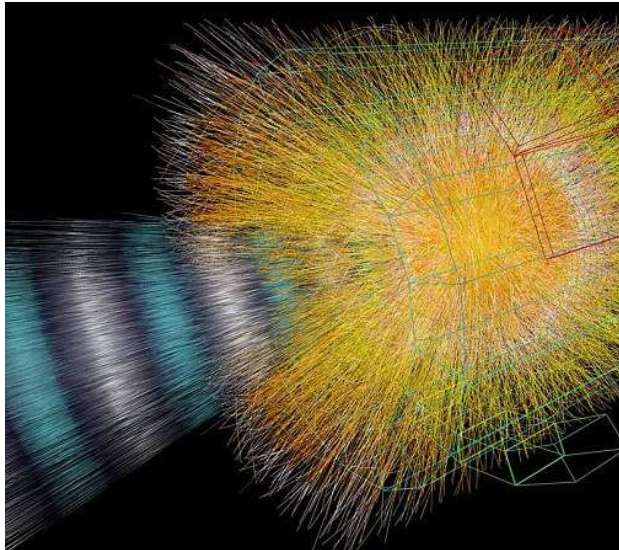
---

- $^{208}\text{Pb} + ^{208}\text{Pb}$  nuclei collide at high energy: 416 nucleons go in,  $\sim 7000$  fundamental particles come out.
- Two protons collide at high energy (7 TeV c.o.m energy for LHC),  $\sim 80$  fundamental particles come out
- The collision converts ordinary matter to pure energy, which then coalesces back out into fundamental particles
- *This is similar to what happened at the Big Bang: start with pure energy, and produce fundamental particles.*
- *So we are recreating conditions of the early universe in these collisions*

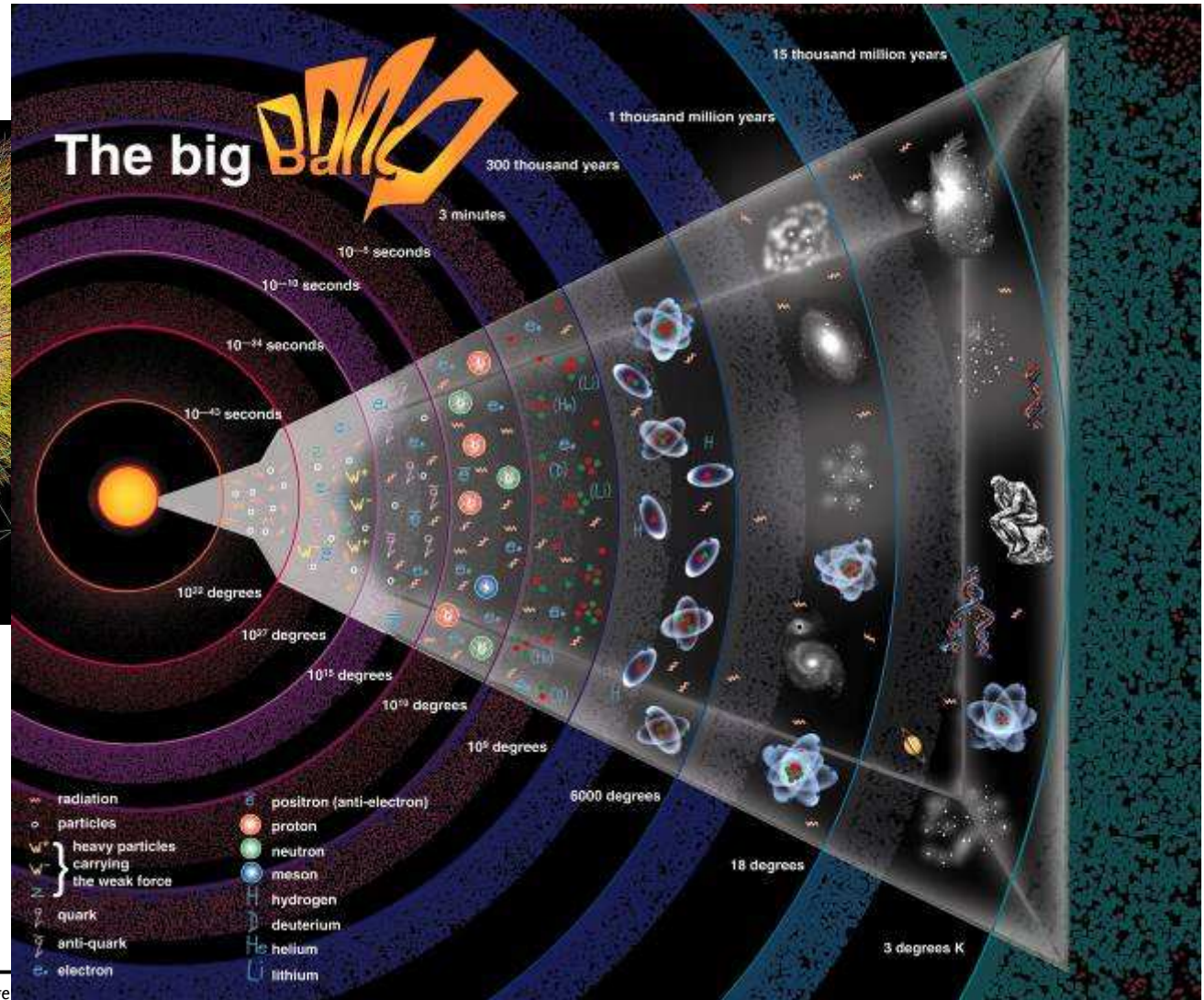


# Why do you want to study fundamental particles ?

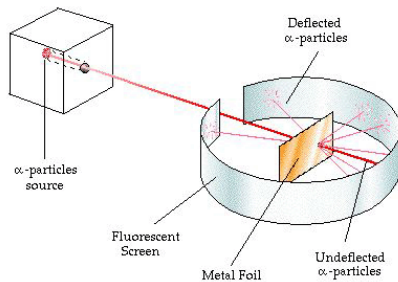
$^{208}\text{Pb} + ^{208}\text{Pb}$  @ 7 TeV



*Recreate conditions  
soon after the BigBang  
& understand how  
matter formed out of  
pure energy*

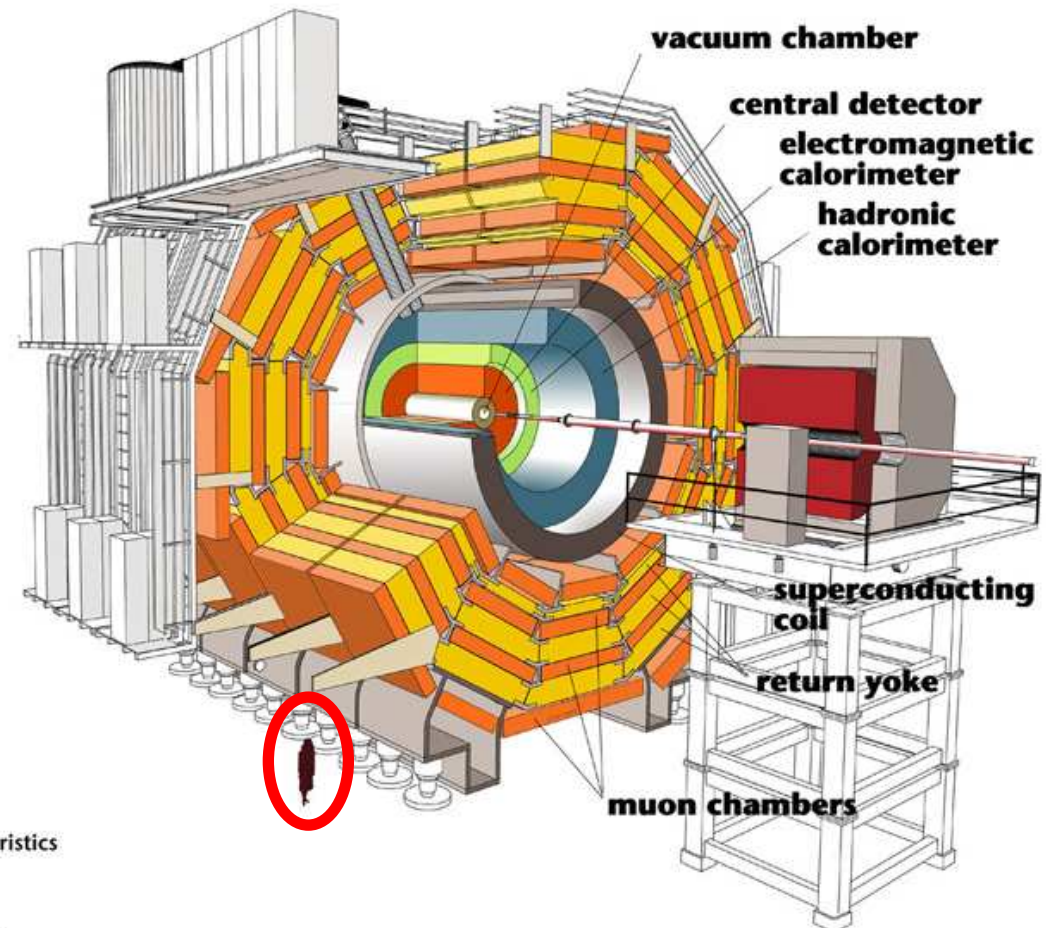


# Why such big experiments?



Rutherford  $\alpha$  experiment

$\sim 1\text{m}$



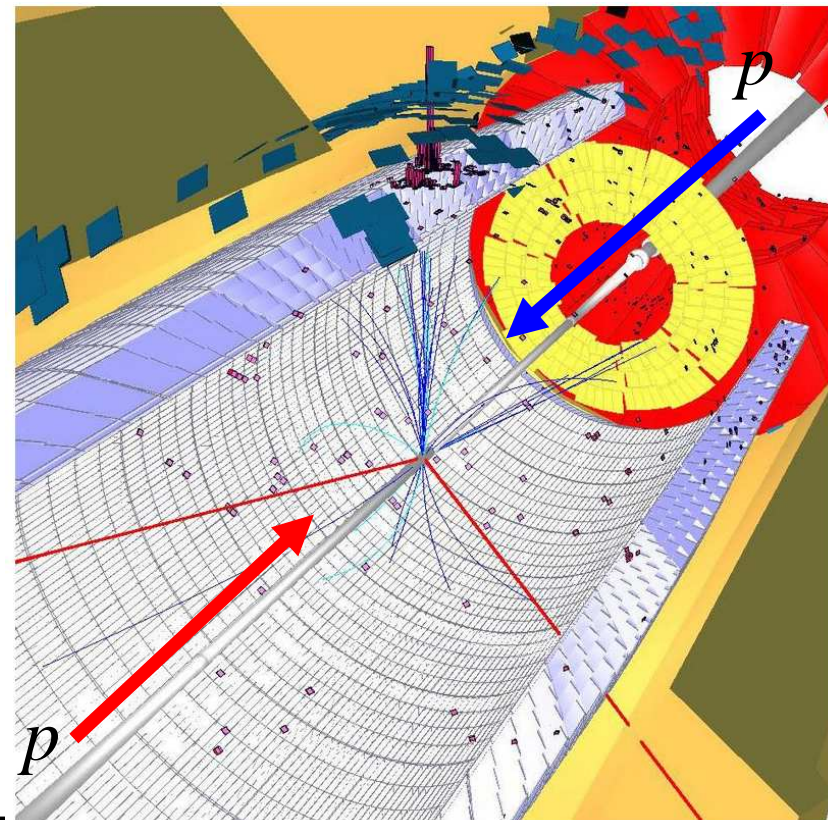
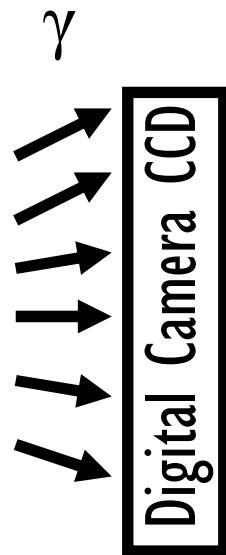
Detector characteristics

Width: 22m  
Diameter: 15m  
Weight: 14'500t



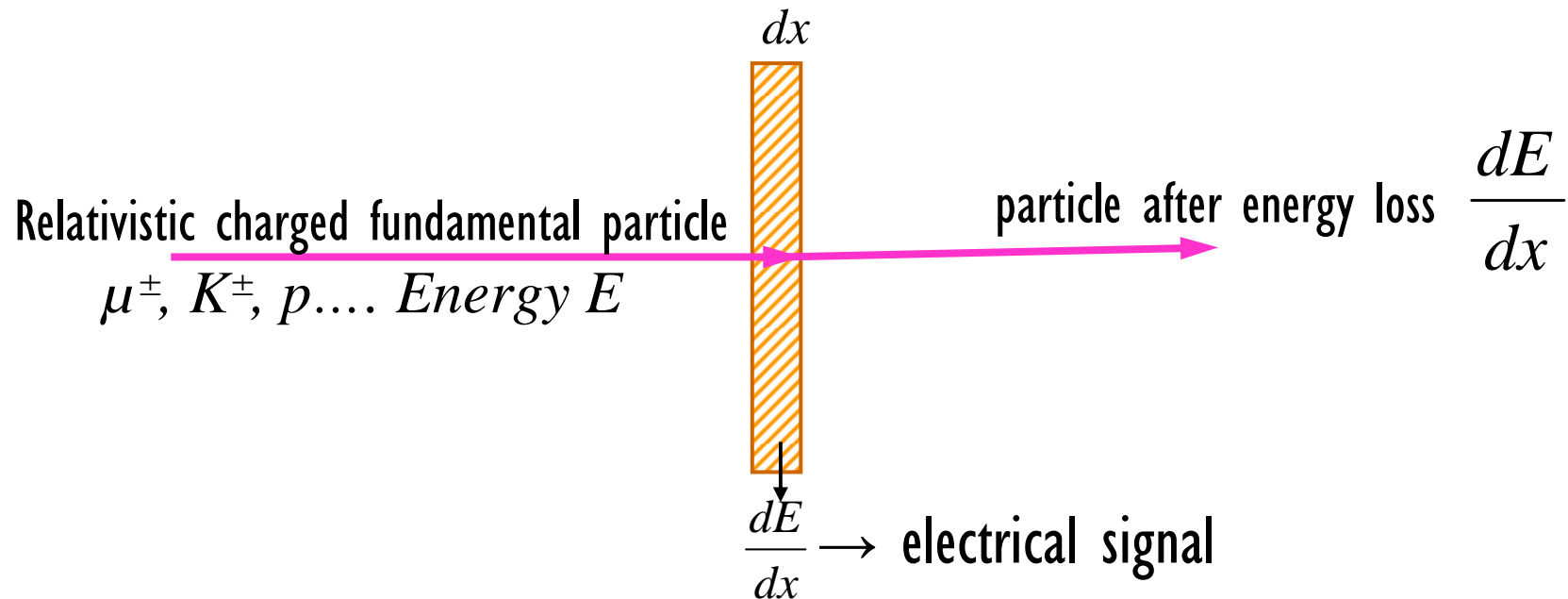
Digital Camera=2D image of photons at low energy

Particle detector=3D image of collision of high energy particles  
with tracking, energy measurements



# Basic principle of particle detection

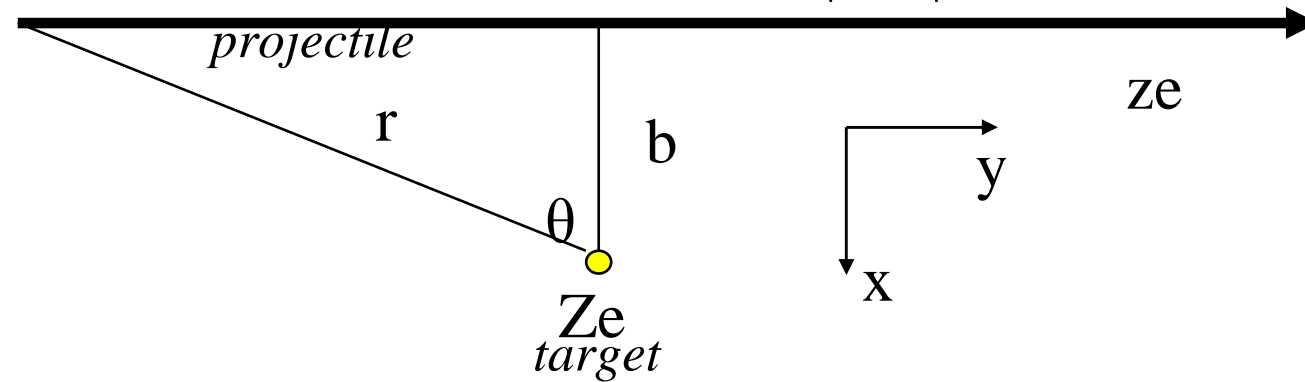
All particles lose energy when passing through a material



There is a fundamental equation that determines  $\frac{dE}{dx}$   
for all charged particles



# Simplified calculation of $\left\langle \frac{dE}{dx} \right\rangle$



Force on projectile

$$F_x = \frac{Zze^2}{4\pi\epsilon_0 r^2} \cos \theta = \frac{Zze^2}{4\pi\epsilon_0 b^2} \cos^3 \theta$$

Force on projectile / target

$$\Delta p = \int_{-\infty}^{\infty} dt F_x = \frac{Zze^2}{2\pi\epsilon_0 \beta c} \frac{1}{b}$$

Energy transferred

$$\Delta E = \frac{\Delta p^2}{2M} = \frac{Z^2 z^2 e^4}{2M (2\pi\epsilon_0)^2 (\beta c)^2} \frac{1}{b^2}$$

target!  $M_{\text{nucleus}} \gg M_{\text{electron}}$   
So scattering is off atomic electrons!

# Simplified calculation of $\left\langle \frac{dE}{dx} \right\rangle$

Energy transferred  $\Delta E$  is a function of the impact parameter 'b':

Integrate over all b's and figure out the average number of electrons as the projectile goes through material of thickness  $\Delta x$

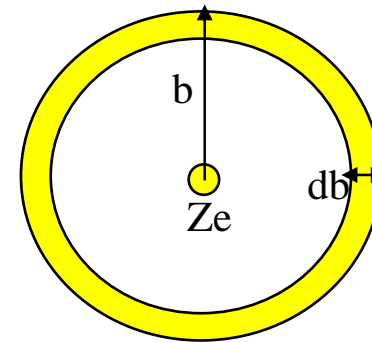
$$\frac{dn}{db} = 2\pi b \times (\text{number of electrons / unit area})$$

$$= 2\pi b \times Z \frac{N_A}{A} \rho \Delta x$$

$$\overline{\Delta E} = \int_{b_{\min}}^{b_{\max}} db \frac{dn}{db} E_e(b) = 2C \frac{m_e c^2}{\beta^2} \frac{Zz^2}{A} \rho \Delta x [\ln b]_{b_{\min}}^{b_{\max}}$$

$$= C \frac{m_e c^2}{\beta^2} \frac{Zz^2}{A} \rho \Delta x [\ln E]_{E_{\min}}^{E_{\max}}$$

with  $C = 2\pi N_A \left( \frac{e^2}{4\pi\epsilon_0 m_e c^2} \right)$



$$E_{\max} = \frac{2\gamma^2 \beta^2 m_e c^2}{1 + 2\gamma \frac{m_e}{M} + \left( \frac{m_e}{M} \right)^2} \approx 2\gamma^2 \beta^2 m_e c^2$$

Maximum energy that a particle with speed  $\beta$  can transfer to a target with mass  $m$  at rest in an elastic collision

$$\left\langle \frac{dE}{dx} \right\rangle$$


---

## Bethe-Bloch equation (1930-33)

Our  
Simplified:

$$\left\langle \frac{dE}{dx} \right\rangle = 2C \frac{m_e c^2}{\beta^2} z^2 \left( \frac{Z}{A} \right) \rho \ln \left( \frac{2\beta^2 \gamma^2 m_e c^2}{I_0} \right)$$

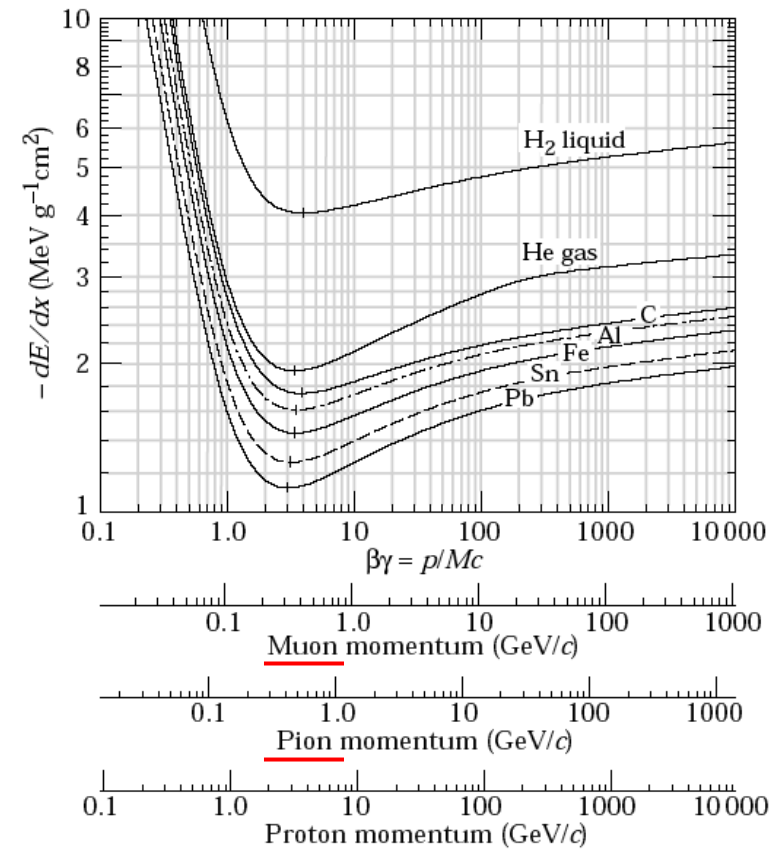
Detailed analysis with quantum mechanics gives:

$$-\frac{dE}{dx} = K z^2 \left( \frac{Z}{A} \right) \frac{1}{\beta^2} \left[ \frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{max}}{I} - \beta^2 - \frac{\delta}{2} \right]$$

# Universal features of Bethe-Bloch Energy Loss

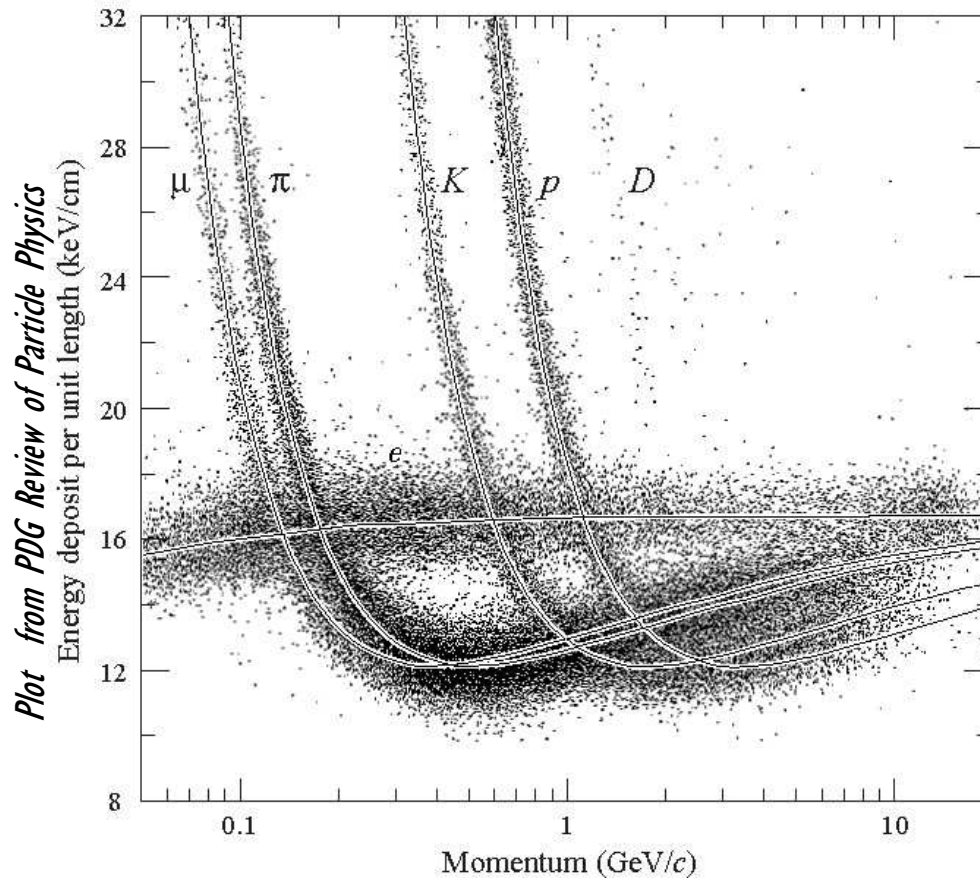
For a relativistic charged particle going through detector material:

1. Shape of energy loss curve is independent of material type & projectile particle ( $\pi^\pm, K^\pm, p..$ )
2. Function of *particle momentum*  
*Z/A of material*  
constants and corrections

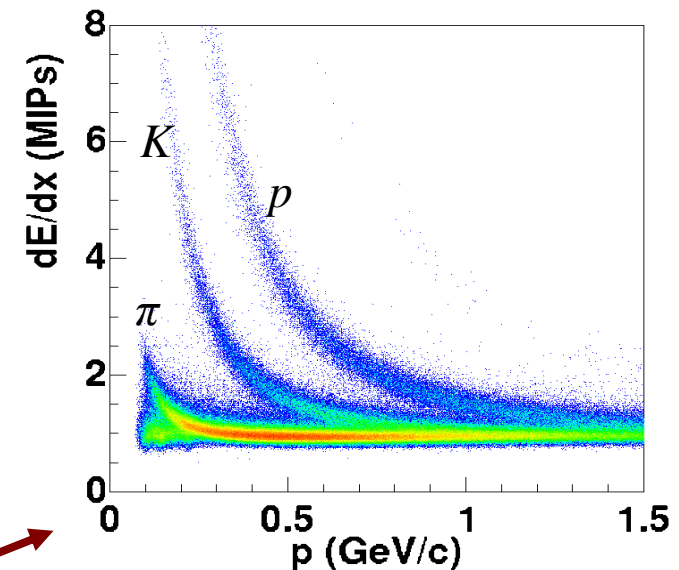




# Why is Bethe-Bloch so important?



Different particles have different  $dE/dx$  in same detector material (due to  $\beta\gamma$  factor)  
→ Particle Identification!



*Plot from Pradeep Sarin, PhD thesis, 300  $\mu$ m thin Silicon pixel detector.*

# What are the fundamental particles we study?

$p$	$P_{11}$	****	$\Delta(1232)$	$P_{33}$	****	$\Sigma^+$	$P_{11}$	****	$\Xi^0$	$P_{11}$	****	$\Lambda_c^+$	****
$n$	$P_{11}$	****	$\Delta(1600)$	$P_{33}$	***	$\Sigma^0$	$P_{11}$	****	$\Xi^-$	$P_{11}$	****	$\Lambda_c(2595)^+$	***
$N(1440)$	$P_{11}$	****	$\Delta(1620)$	$S_{31}$	****	$\Sigma^-$	$P_{11}$	****	$\Xi(1530)$	$P_{13}$	****	$\Lambda_c(2625)^+$	***
$N(1520)$	$D_{13}$	****	$\Delta(1700)$	$D_{33}$	****	$\Sigma(1385)$	$P_{13}$	****	$\Xi(1620)$	*	*	$\Lambda_c(2765)^+$	*
$N(1535)$	$S_{11}$	****	$\Delta(1750)$	$P_{31}$	*	$\Sigma(1480)$	*	*	$\Xi(1690)$	***	***	$\Lambda_c(2880)^+$	***
$N(1650)$	$S_{11}$	****	$\Delta(1900)$	$S_{31}$	**	$\Sigma(1560)$	**	*	$\Xi(1820)$	$D_{13}$	***	$\Lambda_c(2940)^+$	***
$N(1675)$	$D_{15}$	****	$\Delta(1905)$	$F_{35}$	****	$\Sigma(1580)$	$D_{13}$	*	$\Xi(1950)$	***	***	$\Sigma_c(2455)$	****
$N(1680)$	$F_{15}$	****	$\Delta(1910)$	$P_{31}$	****	$\Sigma(1620)$	$S_{11}$	**	$\Xi(2030)$	***	***	$\Sigma_c(2520)$	***
$N(1700)$	$D_{13}$	***	$\Delta(1920)$	$P_{33}$	***	$\Sigma(1660)$	$P_{11}$	***	$\Xi(2120)$	*	*	$\Sigma_c(2800)$	***
$N(1710)$	$P_{11}$	***	$\Delta(1930)$	$D_{35}$	***	$\Sigma(1670)$	$D_{13}$	****	$\Xi(2250)$	**	**	$\Xi_c^+$	***
$N(1720)$	$P_{13}$	****	$\Delta(1940)$	$D_{33}$	*	$\Sigma(1690)$	**	**	$\Xi(2370)$	**	**	$\Xi_c^0$	***
$N(1900)$	$P_{13}$	**	$\Delta(1950)$	$F_{37}$	****	$\Sigma(1750)$	$S_{11}$	***	$\Xi(2500)$	*	*	$\Xi_c^{++}$	***
$N(1990)$	$F_{17}$	**	$\Delta(2000)$	$F_{35}$	**	$\Sigma(1770)$	$P_{11}$	*				$\Xi_c^0$	***
$N(2000)$	$F_{15}$	**	$\Delta(2150)$	$S_{31}$	*	$\Sigma(1775)$	$D_{15}$	****	$\Omega^-$	****	****	$\Xi_c(2645)$	***
$N(2080)$	$D_{13}$	**	$\Delta(2200)$	$G_{37}$	*	$\Sigma(1840)$	$P_{13}$	*	$\Omega(2250)^-$	***	***	$\Xi_c(2790)$	***
$N(2090)$	$S_{11}$	*	$\Delta(2300)$	$H_{39}$	**	$\Sigma(1880)$	$P_{11}$	**	$\Omega(2380)^-$	**	**	$\Xi_c(2815)$	***
$N(2100)$	$P_{11}$	*	$\Delta(2350)$	$D_{35}$	*	$\Sigma(1915)$	$F_{15}$	****	$\Omega(2470)^-$	**	**	$\Xi_c(2930)$	*
$N(2190)$	$G_{17}$	****	$\Delta(2390)$	$F_{37}$	*	$\Sigma(1940)$	$D_{13}$	***				$\Xi_c(2980)$	***
$N(2200)$	$D_{15}$	**	$\Delta(2400)$	$G_{39}$	**	$\Sigma(2000)$	$S_{11}$	*				$\Xi_c(3055)$	**
$N(2220)$	$H_{19}$	****	$\Delta(2420)$	$H_{3,11}$	****	$\Sigma(2030)$	$F_{17}$	****				$\Xi_c(3080)$	***
$N(2250)$	$G_{19}$	****	$\Delta(2750)$	$I_{3,13}$	**	$\Sigma(2070)$	$F_{15}$	*				$\Xi_c(3123)$	*
$N(2600)$	$I_{1,11}$	***	$\Delta(2950)$	$K_{3,15}$	**	$\Sigma(2080)$	$P_{13}$	**				$\Omega_c^0$	***
$N(2700)$	$K_{1,13}$	**				$\Sigma(2100)$	$G_{17}$	*				$\Omega_c(2770)^0$	***
			$\Lambda$	$P_{01}$	****	$\Sigma(2250)$		***					
			$\Lambda(1405)$	$S_{01}$	****	$\Sigma(2455)$		**				$\Xi_{cc}^+$	*
			$\Lambda(1520)$	$D_{03}$	****	$\Sigma(2620)$		**					
			$\Lambda(1600)$	$P_{01}$	***	$\Sigma(3000)$		*				$\Lambda_b^0$	***
			$\Lambda(1670)$	$S_{01}$	****	$\Sigma(3170)$		*				$\Sigma_b$	***
			$\Lambda(1690)$	$D_{03}$	****							$\Sigma_b^*$	***
			$\Lambda(1800)$	$S_{01}$	***							$\Xi_b^0, \Xi_b^-$	***
			$\Lambda(1810)$	$P_{01}$	***							$\Omega_b^-$	***
			$\Lambda(1820)$	$F_{05}$	****								
			$\Lambda(1830)$	$D_{05}$	****								
			$\Lambda(1890)$	$P_{03}$	****								
			$\Lambda(2000)$	*									
			$\Lambda(2020)$	$F_{07}$	*								
			$\Lambda(2100)$	$G_{07}$	****								
			$\Lambda(2110)$	$F_{05}$	***								
			$\Lambda(2220)$	$\Sigma$	*								

*Particle Data Book lists ~ 200 known  
“stable” fundamental particles “Particle Zoo”*

# Most of those particles don't travel far

---

Distance a particle travels (and can be detected before it decays) depends on its energy.

Example 1:  $\mu \rightarrow e^- + \bar{\nu}_e + \nu_\mu$  lifetime  $\tau = 2.2 \times 10^{-6} s$

- Cosmic ray  $\mu$ 's produced by cosmic rays colliding with  $p, He, Li$  in upper atmosphere: altitude  $\sim 10$  km.
- In its rest frame  $\tau \sim 2.2 \cdot 10^{-6} s$ . At speed  $\sim c$  expect  $\mu$  to travel  $\sim c\tau \sim 660$  m before decaying to  $e^-$ .  $m_\mu c^2 = 105$  MeV
- **But**  $E_\mu \sim 2$  GeV so it has a Lorentz boost of  $E_\mu = \gamma (m_\mu c^2)$  so  $\gamma \sim 20$  gives mean range of  $\mu$  before decay  $s = c\gamma\tau \sim 12$  km

Example 2:

$\pi$  has shorter lifetime  $2.6 \cdot 10^{-8} s$      $\rho_0$  lifetime  $\sim 6 \cdot 10^{-24} s$

# Reducing the Zoo membership in a $p+p$ collider @ 7 TeV

---

1. From the  $\sim 200$  fundamental particles listed by the PDG, only **27** have a  $c\gamma\tau > \sim 1\mu m$  so they can be seen as 'tracks' in a detector.
2. **13** of these have a  $c\gamma\tau < 500\mu m$ , *i.e. very short tracks* that must be measured indirectly with high precision vertex detectors.
3. Of the **14** remaining particles only **8** typically have energies high enough to have  $c\gamma\tau > \text{few tens of meters}$  to make it all the way through most detectors:

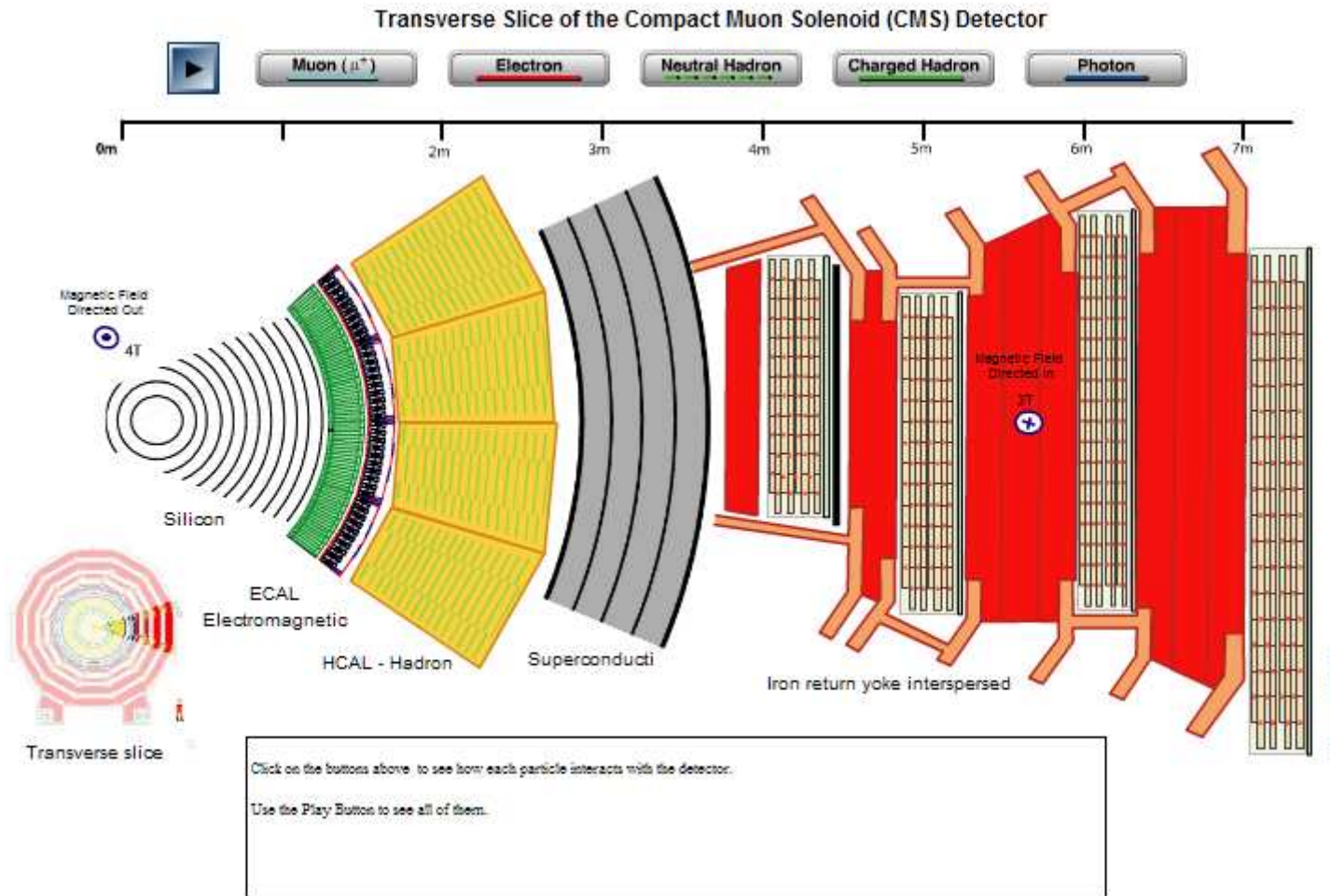
$$e^{\pm}, \mu^{\pm}, \gamma^0, \pi^{\pm}, K^{\pm}, K^0, p^{\pm}, n$$

*These are the common 'bread and butter' of particle physics.*



# Particle detection at work

## A slice through the CMS detector



Derived from CMS Detector Slice from CERN

# What are the measurements we make?

---

- **Energy  $E$**

through energy deposition and/or calorimetry

- **Momentum  $p$**

make charged particles bend in a magnetic field

- $(E^2 - p^2) = m^2$  uniquely identifies particles

- Can also measure spin, polarization in specialized experiments.

# Summary

---

In this lecture we have looked at:

- Why we do high energy Particle & Nuclear physics
- What are the common fundamental particles we deal with
- Principle of particle detection through energy loss